



Is Polarisation the Key to Understanding Magnetar Emission?

Mode Conversion in a Magnetar Atmosphere

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X-ray polarization in magnetar atmospheres – effects of mode conversion

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tertiary Imaging
– served four
– take

ABSTRACT
Magnetars, the most strongly magnetized neutron stars, are an X-ray Polarimetry Explorer (XPPE), the first satellite devoted to magnetars to date. A proper interpretation of the polarization measurements requires a proper account of the effects of the magnetized vacuum or plasma on the propagation of the radiation. In this paper we study the effects of mode conversion in the vacuum resonance of magnetar atmospheres under varying conditions of magnetic field strength and geometry. We show that mode conversion can significantly affect the polarization signal, especially in the case of partial adiabatic mode conversion. Complete mode conversion at the vacuum resonance (O) or partial adiabatic mode conversion (X) to the ordinary (O) or extraordinary (X) modes is observed for magnetic field strengths of $B \approx 10^{13}$ G or partial adiabatic mode conversion with a standard plasma at ≈ 0.5 keV for a magnetic field of $B = 3 \times 10^{13}$ G. We discuss the implications for the interpretation of the polarization signal at infinity for different degrees when compared with a standard plasma at ≈ 0.5 keV. We also discuss the expected polarization signal at infinity for different degrees when compared with a standard plasma at ≈ 0.5 keV. We also discuss the observability of quantum electrodynamic (QED) effects in the magnetar atmosphere. Demonstration of a Soft X-ray Polarimeter (SXP) polarization – radiative transfer

Key words: polarization – radiative transfer

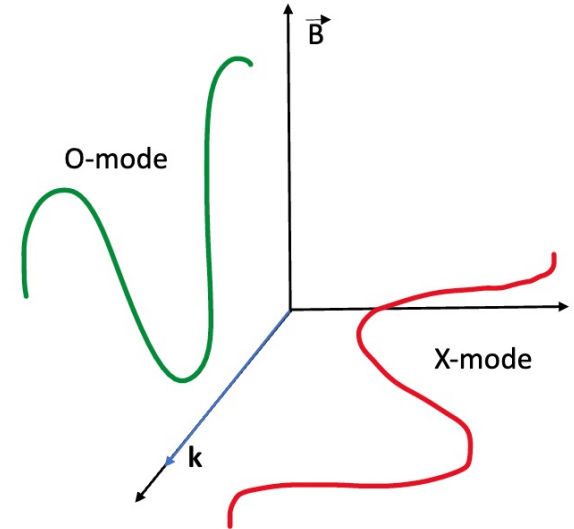
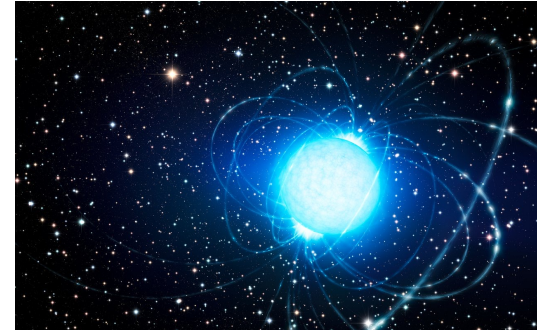
Magnetars and Polarisation

Magnetars:

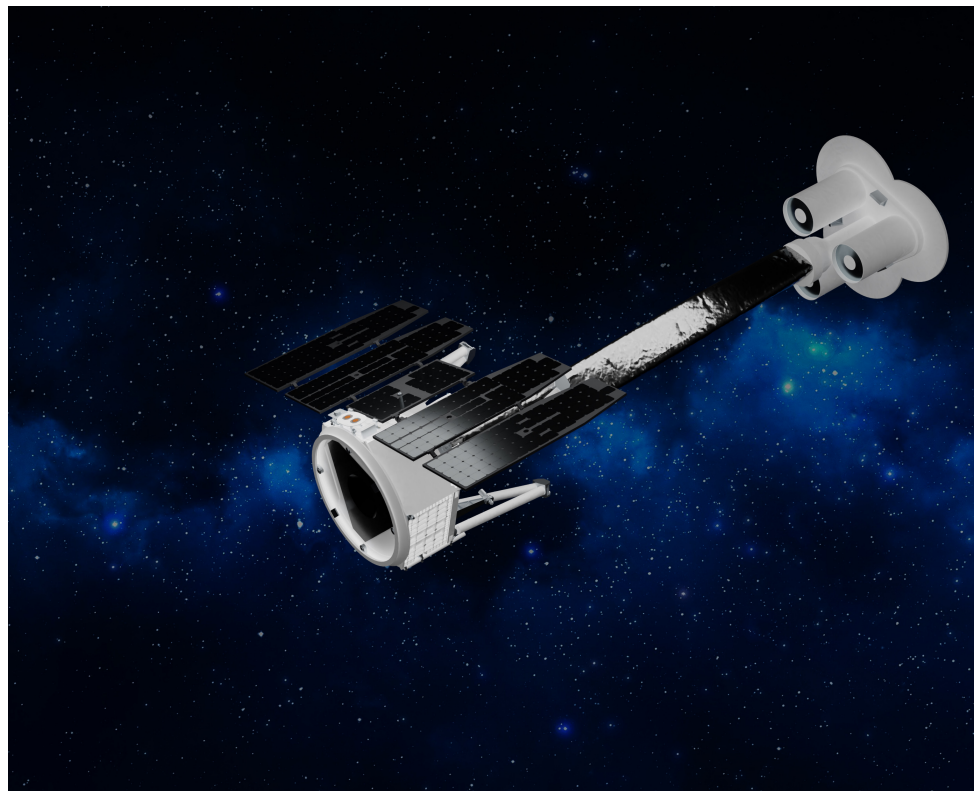
- NS with magnetic field above $10^{13} G$
- Detected in X-ray range

Polarisation:

- Anisotropic atmospheric medium
- Electromagnetic waves polarised in two normal modes
- X-mode has reduced refractive index
- Polarisation degree
- Polarisation angle



- Exploit polarisation state of light to investigate X-ray production
- 2-8 keV
- Four magnetars to date:
 - 4U 0142+61
 - IRXS J1708
 - SGR 1806
 - IE 2259+586



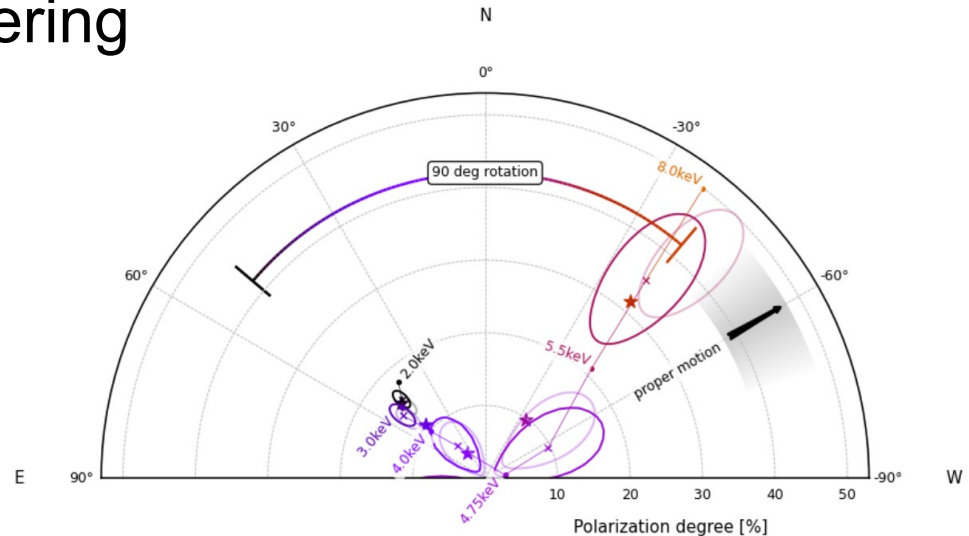
$$B \approx 1.3 \times 10^{14} \text{ G}$$

Taverna et al 2022

- Resonant Compton Scattering
- Condensed surface

Lai 2023

- Partial mode conversion



$$B \approx 4 - 5 \times 10^{14} \text{ G}$$

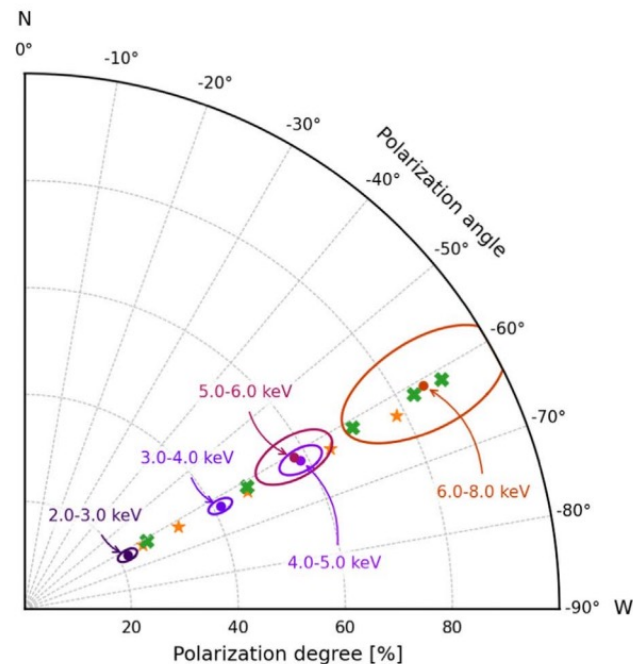
Zane et al 2023

High polarisation

➤ Standard atmosphere emission

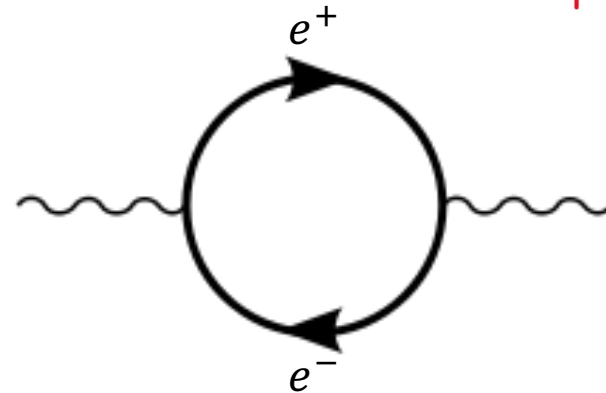
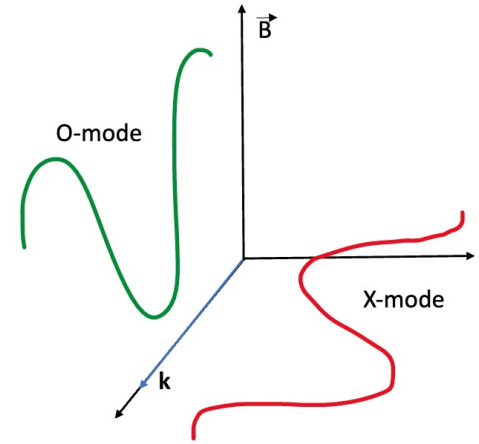
Low polarisation

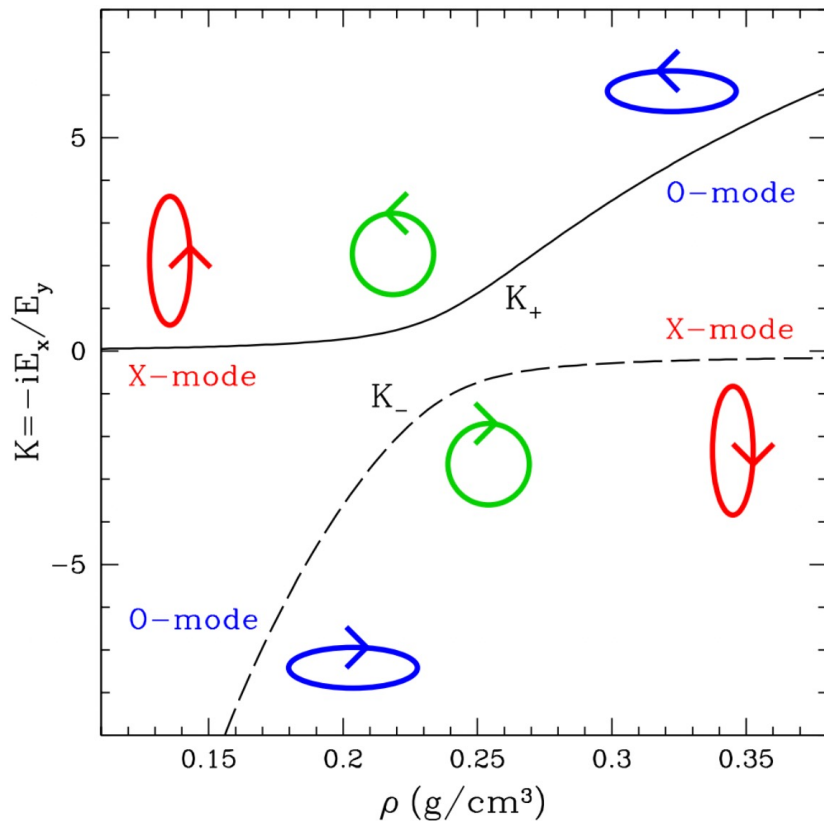
➤ Warm condensed region



Mode Conversion

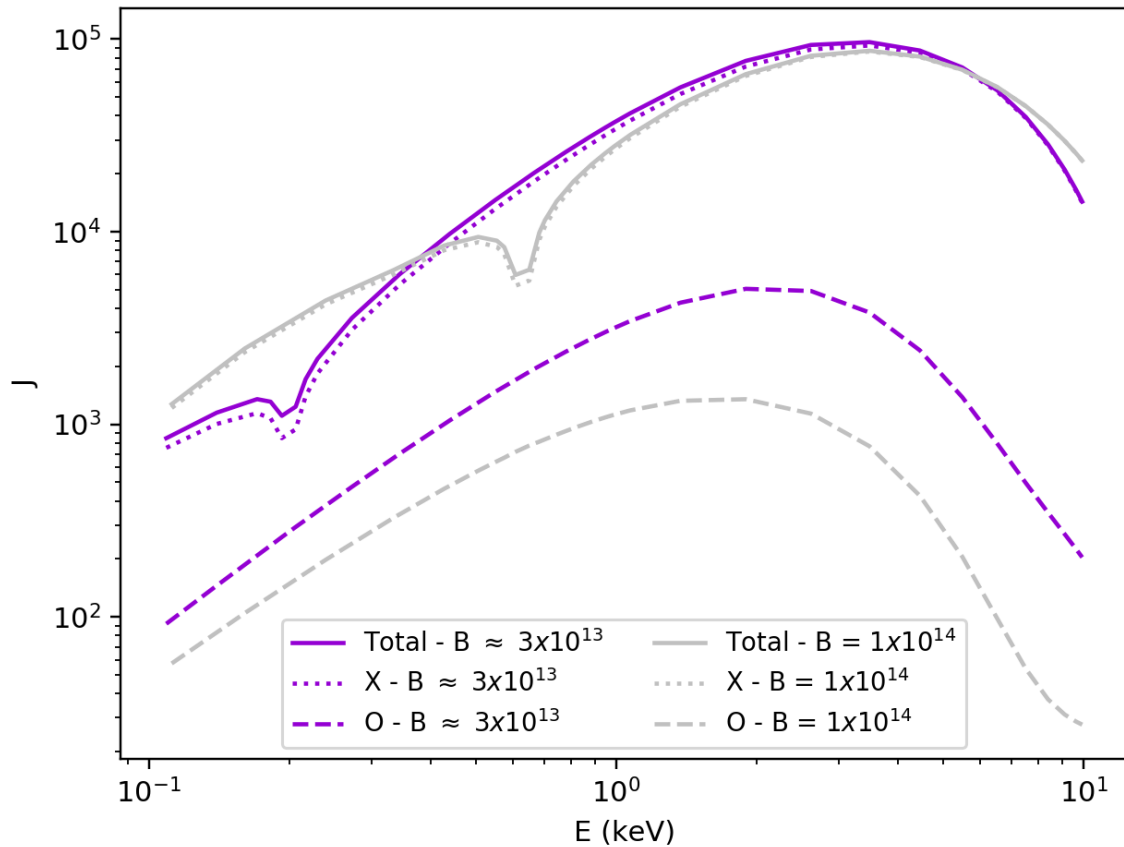
- Two polarisation modes
- High magnetic fields
- Vacuum birefringence (QED)
- Both plasma and vacuum contribute to dielectric tensor
- “Vacuum resonance”



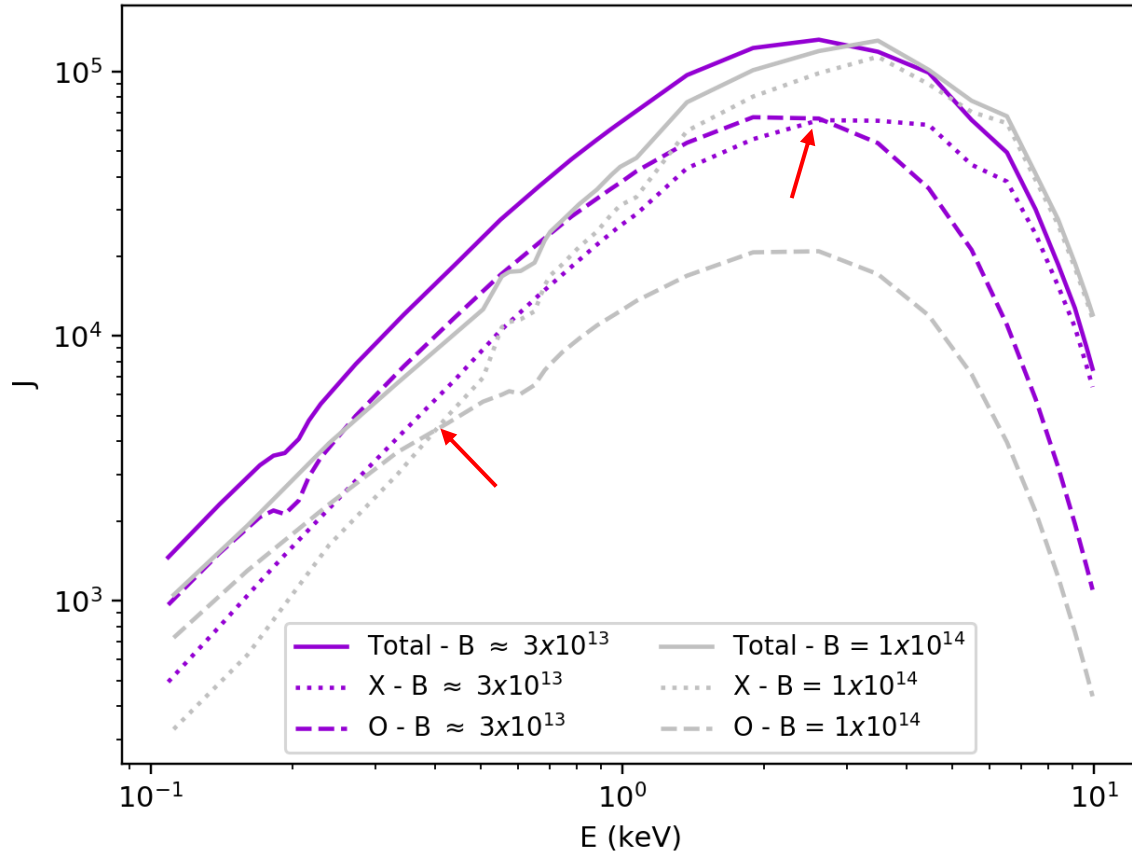


- No mode conversion
- Complete mode conversion
- Partial mode conversion
 - Adiabatic mode conversion

Standard Atmospheric Emission



Complete Mode Conversion



Probability of Mode Conversion

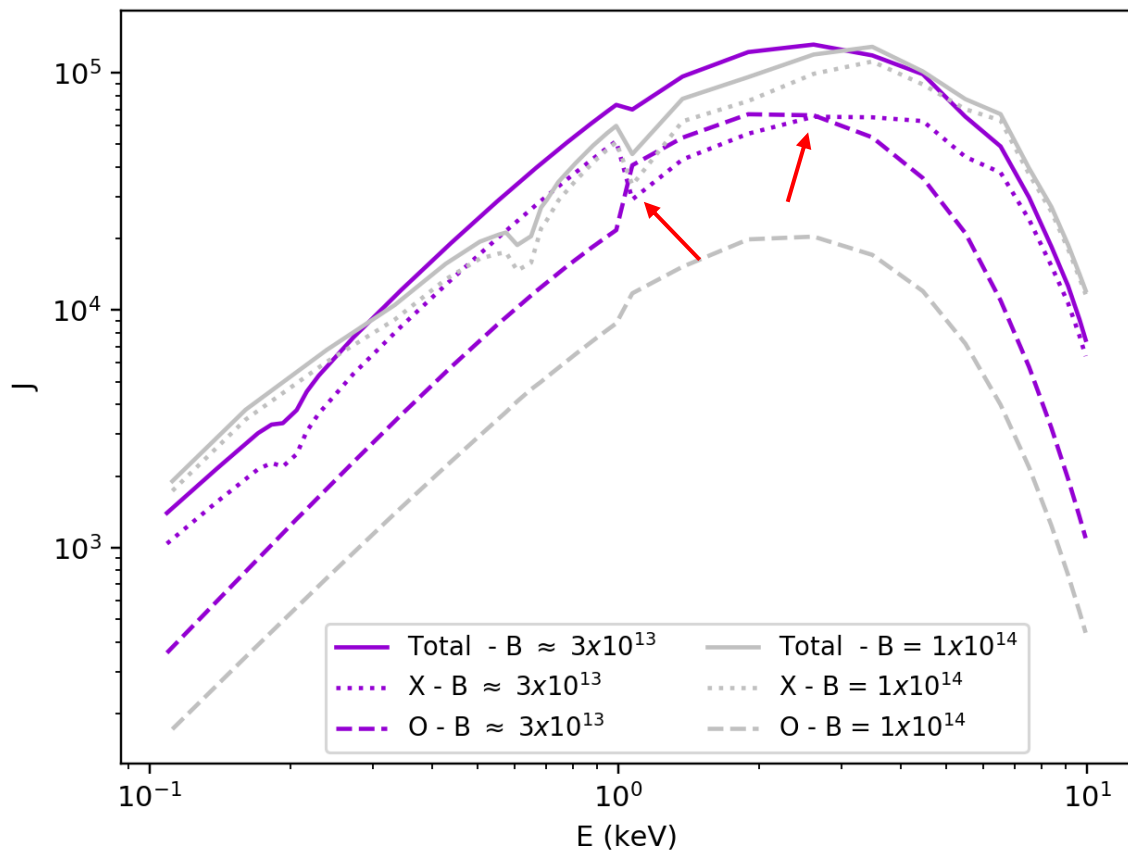
$$1 - \exp \left[-\frac{\pi}{2} \left(\frac{E}{E_{ad}} \right)^3 \right],$$

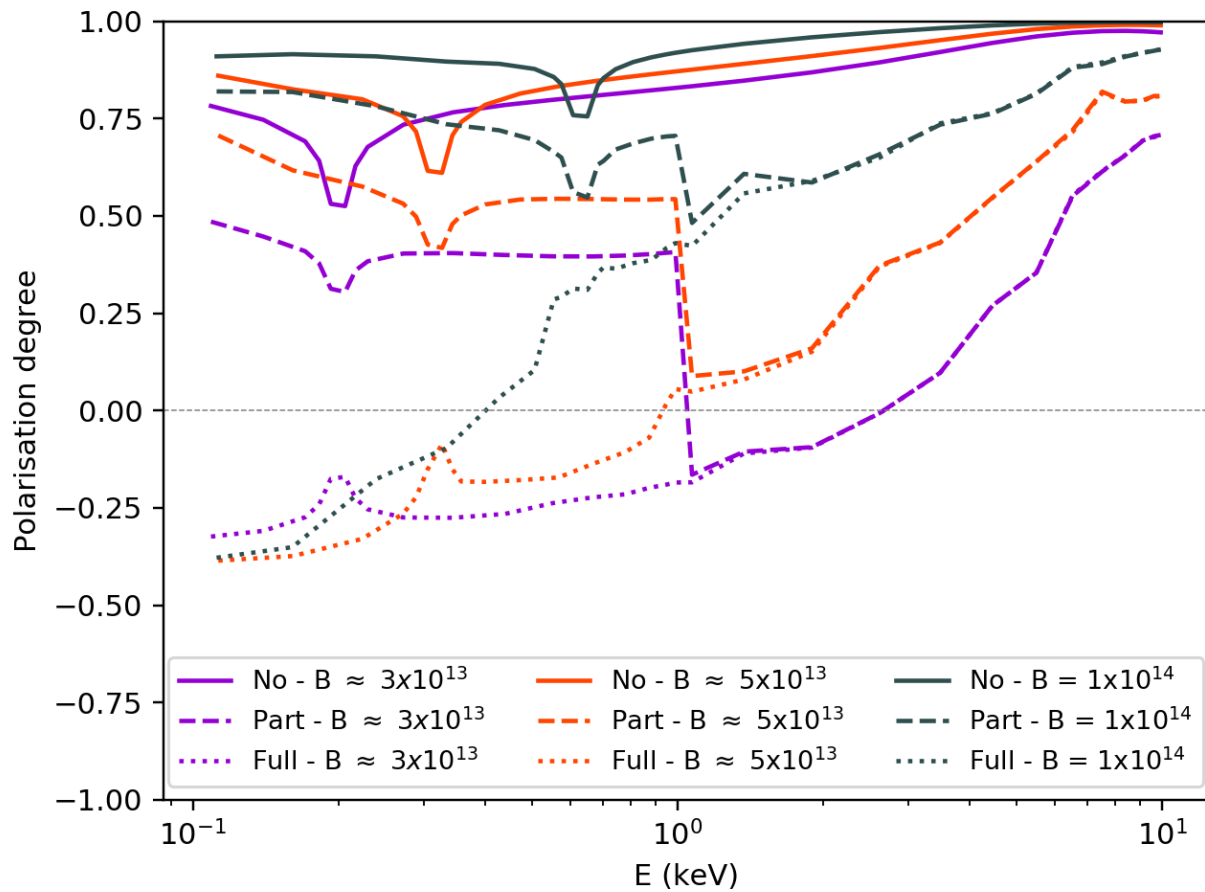
$$E_{ad} = 2.52 (f \tan \theta_B)^{\frac{2}{3}} \left(\frac{1 \text{ cm}}{H_p} \right)^{\frac{1}{3}} \text{ keV}$$

Every photon with $P > P_{con}$:
mode conversion occurs

$$\begin{aligned} f &= 1 \\ \tan \theta_B &= 1 \\ H_p &= 1 \end{aligned}$$

Partial Mode Conversion ($P_{con} = 0.1$)

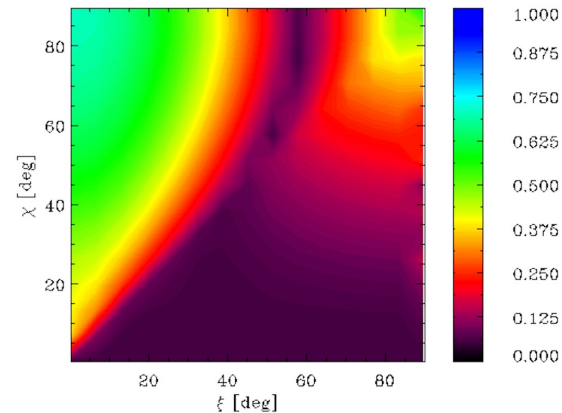
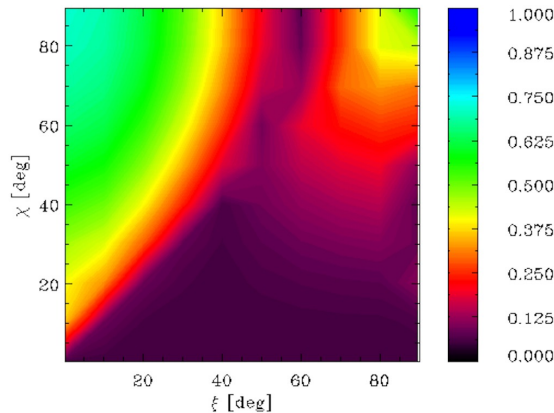
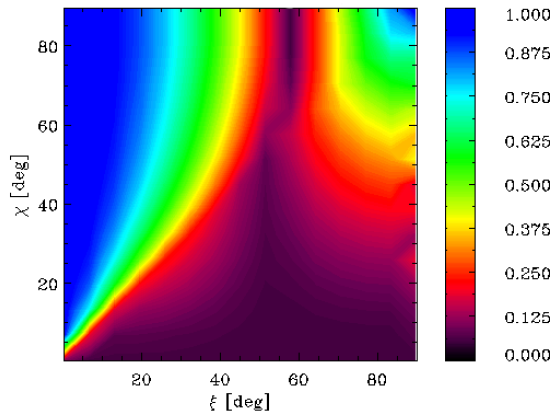




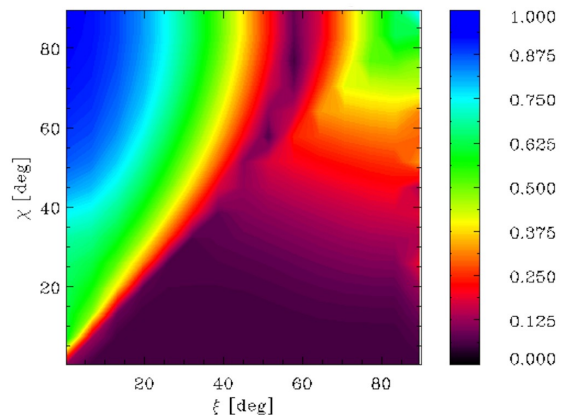
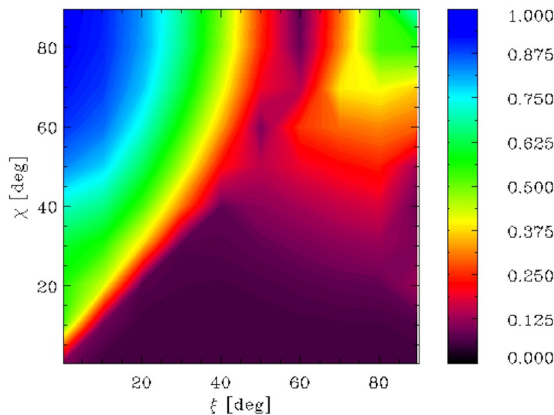
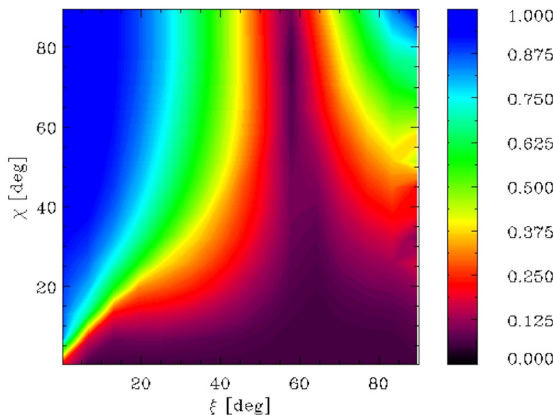
$$\text{PD}_{\text{em}} = \frac{J_{\nu}^{\text{X}} - J_{\nu}^{\text{O}}}{J_{\nu}^{\text{X}} + J_{\nu}^{\text{O}}}$$

$$B_{pol} = 1 \times 10^{14}$$

2-4 keV



6-8 keV



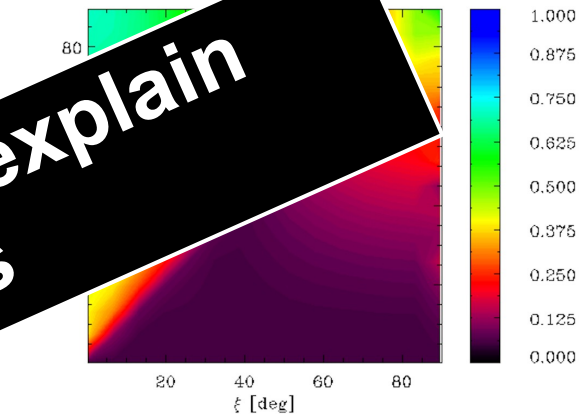
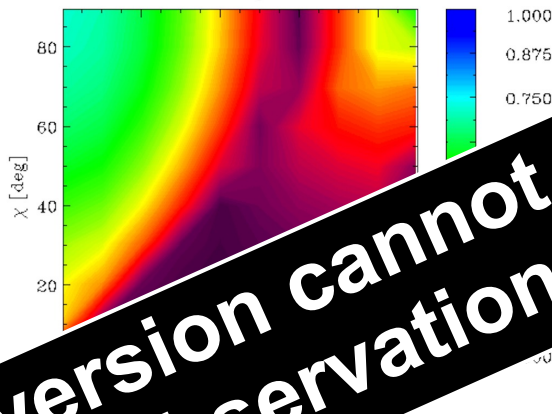
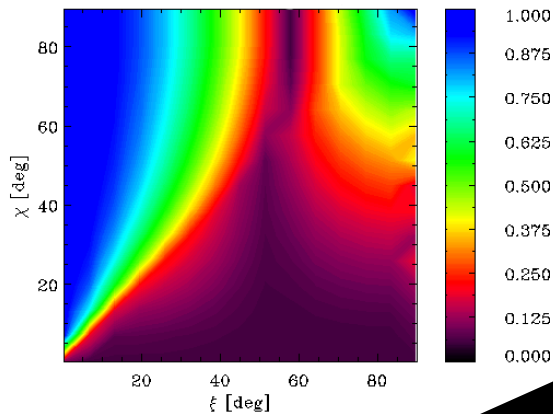
Standard

Complete

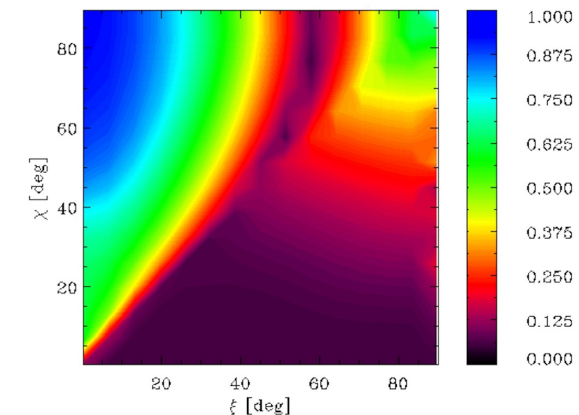
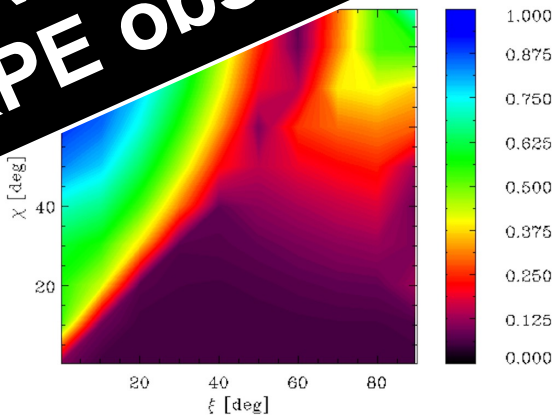
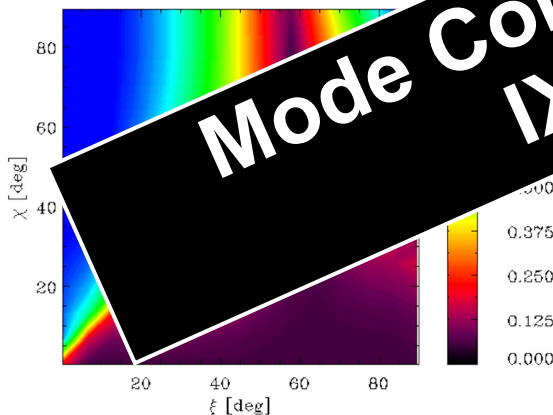
Partial

$$B_{pol} = 1 \times 10^{14}$$

2-4 keV



6-8 keV



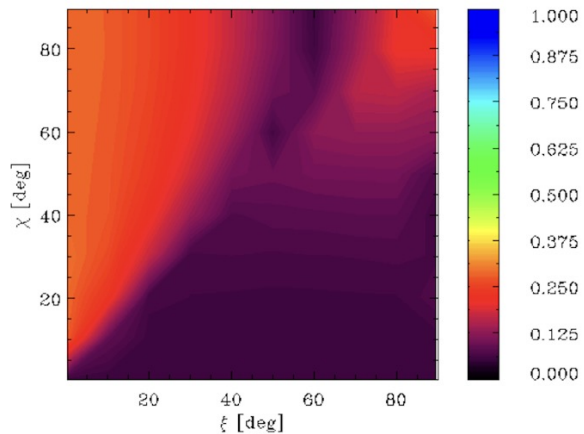
Standard

Complete

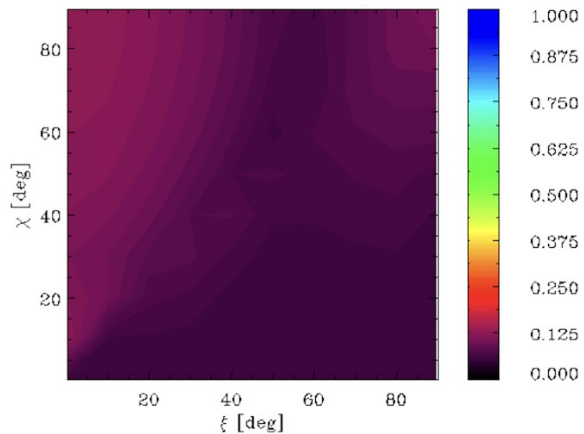
Partial

Mode Conversion cannot explain IXPE observations

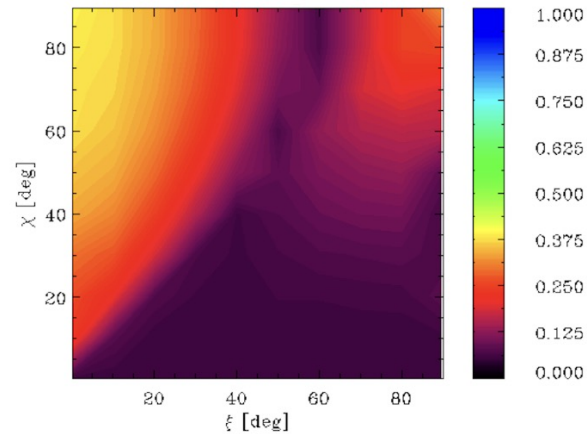
Complete Mode Conversion



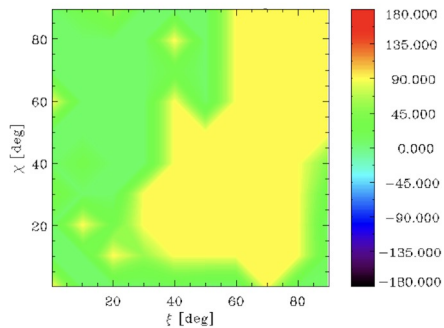
0.1-0.5 keV



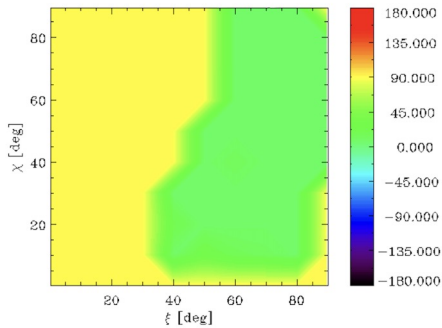
0.5-1 keV



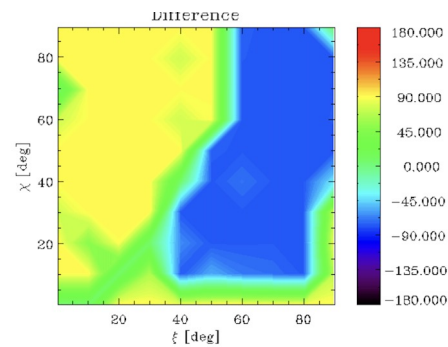
1-2 keV



0.1-0.5 keV

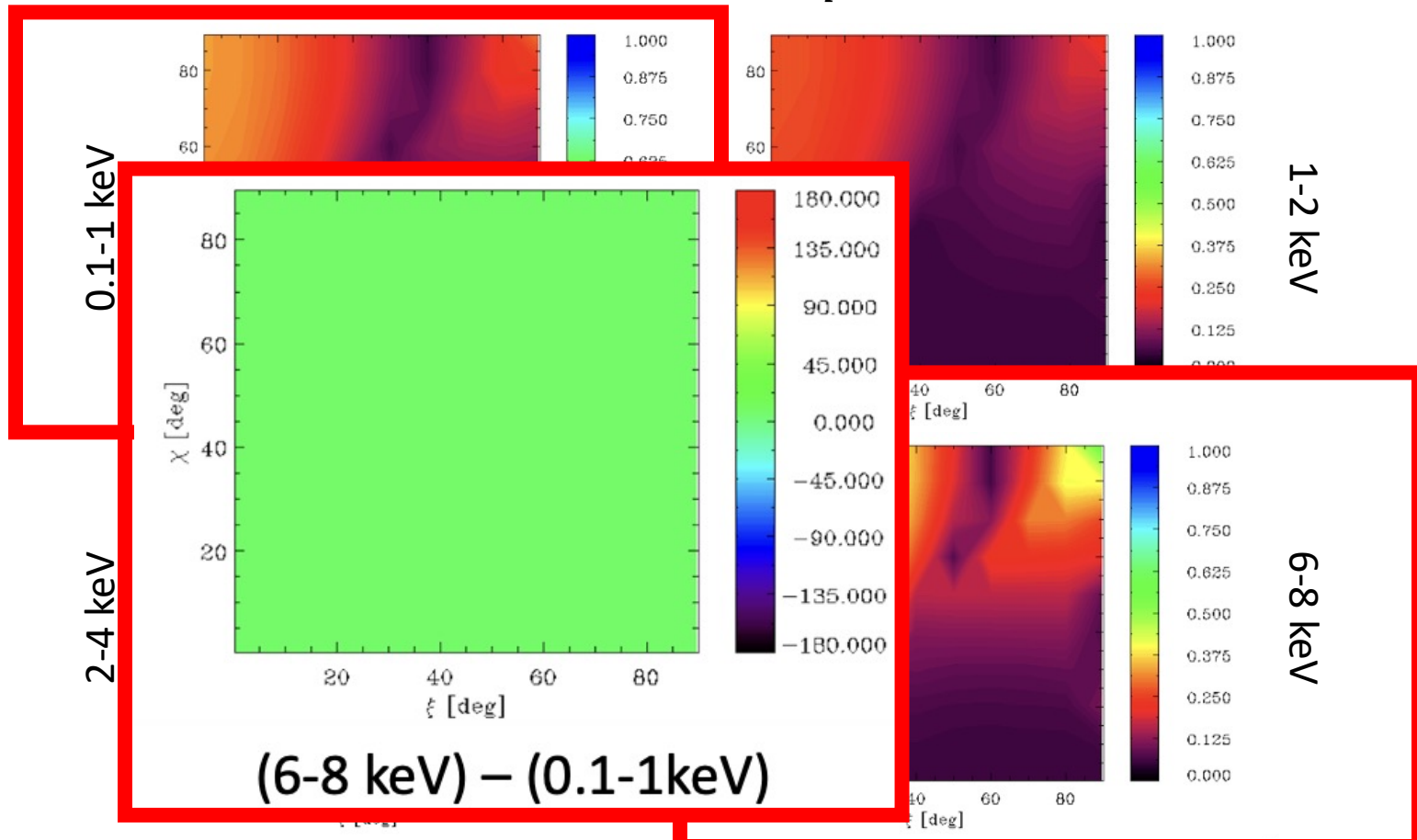


1-2 keV



Difference

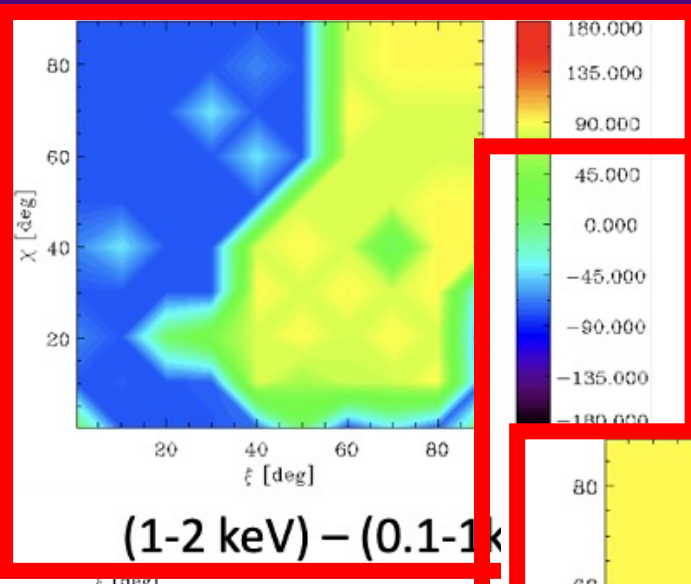
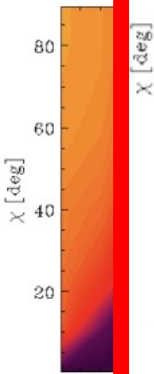
Partial Mode Conversion ($B_{pol} = 5 \times 10^{13}$ G)



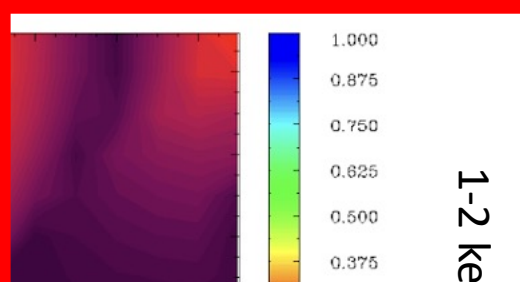
Par

$\mu = 5 \times 10^{13}$ G)

0.1-1 keV

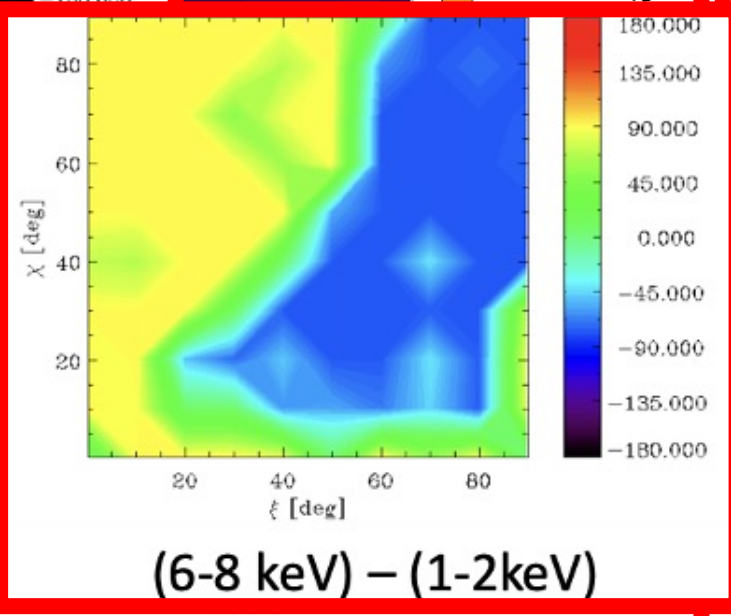
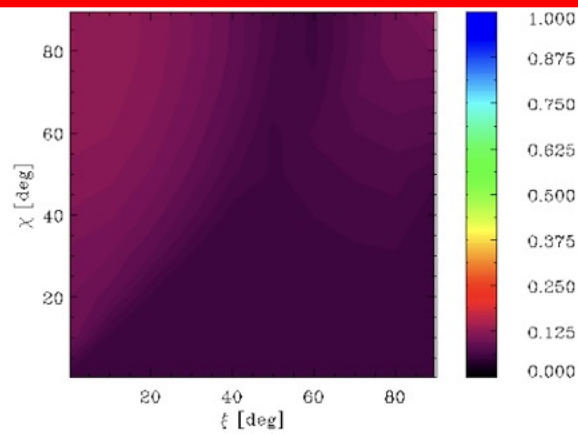


(1-2 keV) - (0.1-1 keV)



1-2 keV

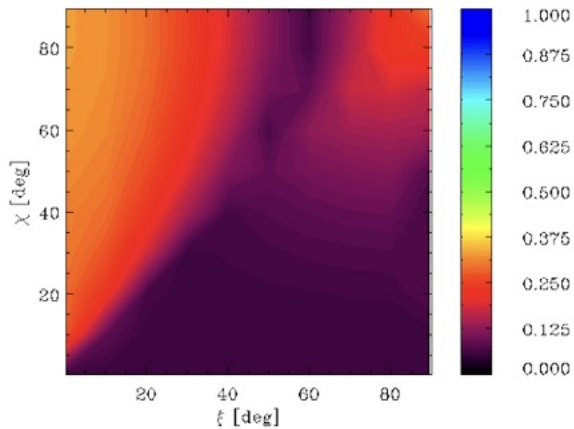
2-4 keV



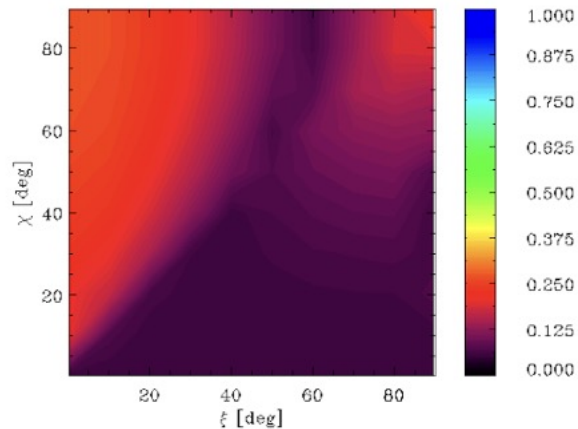
(6-8 keV) - (1-2 keV)

Partial Mode Conversion ($B_{pol} = 5 \times 10^{13}$ G)

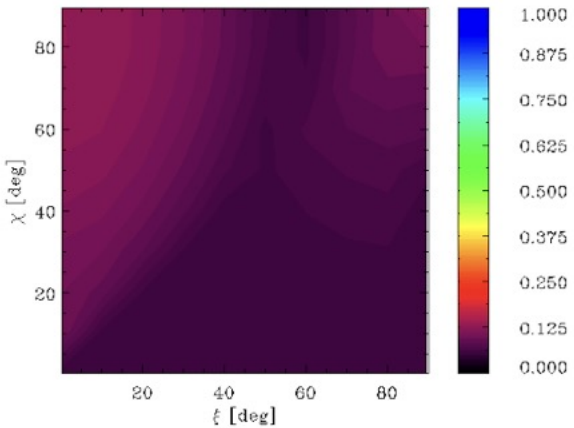
0.1-1 keV



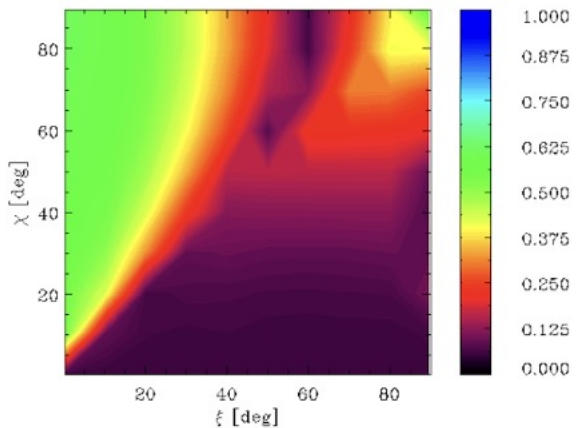
1-2 keV



2-4 keV



6-8 keV



Conclusion

- Complete MC – switch in dominant polarisation mode ($E \uparrow$ as $B \downarrow$)
- Partial MC – reduced polarisation – two switches at low enough B
 - In IXPE – **XDINS** ?
 - **REDSOX** (0.2-0.8 keV) ?



Future and Ongoing work

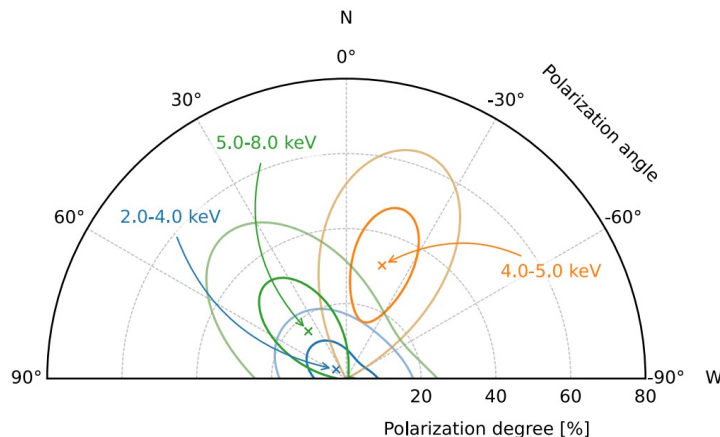
- Particle bombardment
 - Partial ionization

Thank you!

$$B \approx 8 \times 10^{14} \text{G}$$

Turolla et al 2023

- Resonant Compton Scattering
- Condensed surface



IE 2259+586

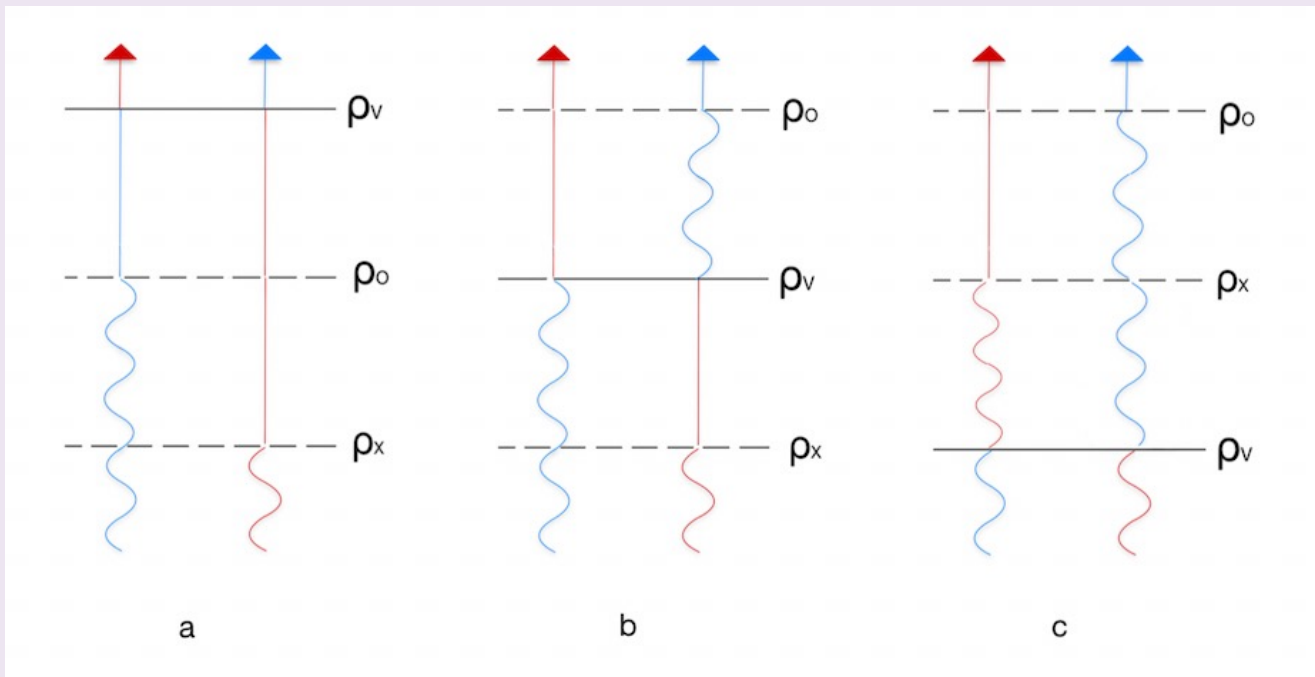
$$B \approx 6 \times 10^{13} \text{G}$$

Or

$$B > 10^{15} \text{G}$$

Heyl et al 2024

- mild, phase-dependent PD
- $\sim 0 - 25 \%$
- a baryon-loaded magnetic loop



$$\rho_V \approx 0.96 Y_e^{-1} \left(\frac{E}{1 \text{keV}} \right)^2 \left(\frac{B}{10^{14}} \right)^2 f(B)^{-2} \text{g cm}^{-3}$$

Ho and Lai, 2003

$$\tau_O = \rho \sigma z$$

$$\sigma_O = \sigma_T \left(\sin^2 \theta + \frac{1}{2} \cos^2 \theta \left[\frac{\epsilon^2}{(\epsilon + \epsilon_B)^2} + \frac{\epsilon^2}{(\epsilon - \epsilon_B)^2} \right] \right)$$

Harding and Lai, 2006