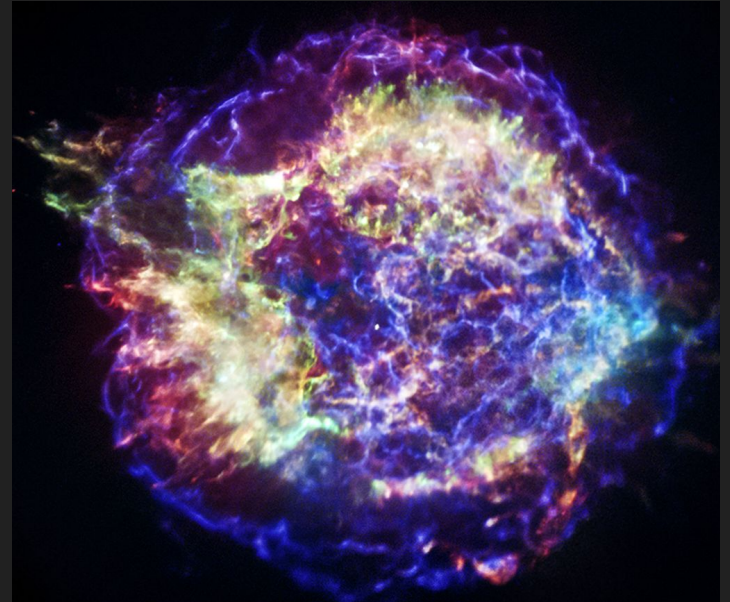


# Neutron star population synthesis: an overview and new results

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in collaboration with Michele Ronchi,  
Celsa Pardo Araujo, and Nanda Rea



Cassiopeia A supernova remnant  
(credit: NASA/CXC/SAO)

# Neutron star population synthesis

- We can estimate the **total number of neutron stars in our Galaxy**:

**CC supernova rate:**  
~ 2 per century

×

**Galaxy age:**  
~ 13.6 billion years

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**NS number:**  
~  $2.8 \times 10^8$

- We only **detect** a very **small fraction** of all neutron stars. Population synthesis bridges this gap focusing on the full population of neutron stars (e.g. Faucher-Giguère & Kaspi 2006, Lorimer et al. 2006, Gullón et al. 2014, Cieřlar et al. 2020):

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model **birth properties** with Monte-Carlo approach



**evolve properties** forward in time

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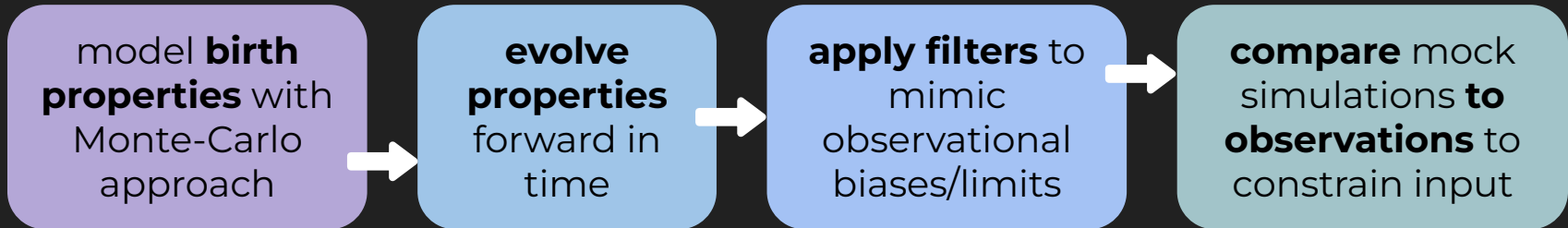
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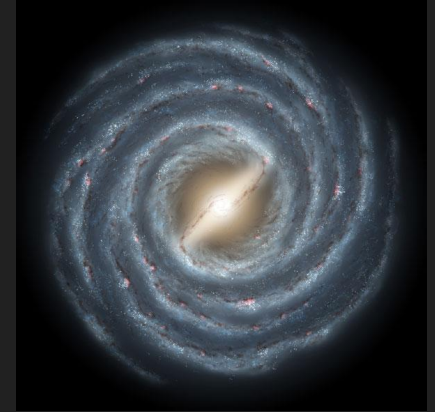
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# Dynamical evolution I

- **Neutron stars are born in star-forming regions**, i.e., in the Galactic disk along the Milky Way's spiral arms, **and receive kicks** during the supernova explosions.
- We make the following assumptions:
  - Electron-density model (Yao et al., 2017) + rigid rotation with  $T = 250$  Myr.
  - **Exponential disk** with scale height  $h_c = 0.18$  kpc (Wainscoat et al., 1992).
  - Single-component **Maxwell kick-velocity distribution** with dispersion  $\sigma_k = 265$  km/s (Hobbs et al., 2005).
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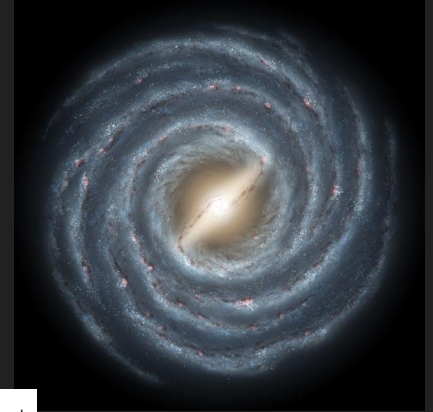
Artistic illustration of the Milky Way (credit: NASA JPL)



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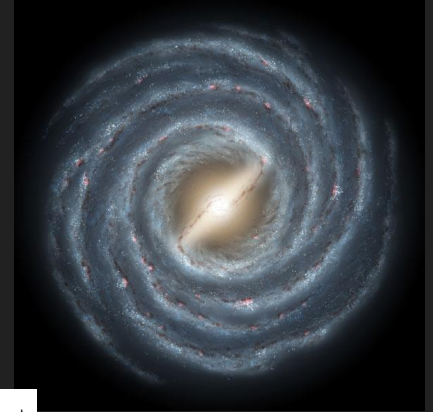
$$\mathcal{P}(z) = \frac{1}{h_c} e^{-\frac{|z|}{h_c}}$$

$$\mathcal{P}(v_k) = \sqrt{\frac{2}{\pi}} \frac{v_k^2}{\sigma_k^3} e^{-\frac{v_k^2}{2\sigma_k^2}}$$

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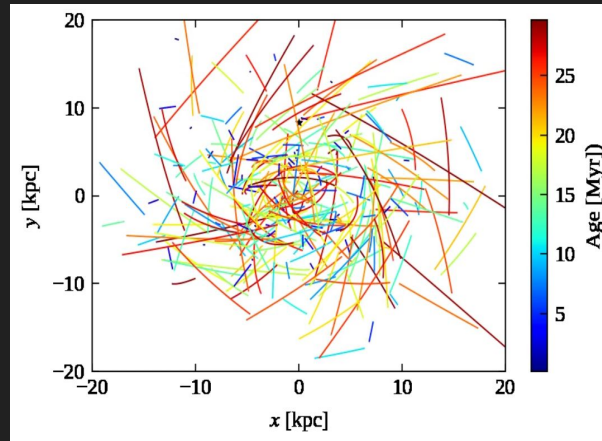
**We use this information to determine pulsar positions and velocities.**

# Dynamical evolution II

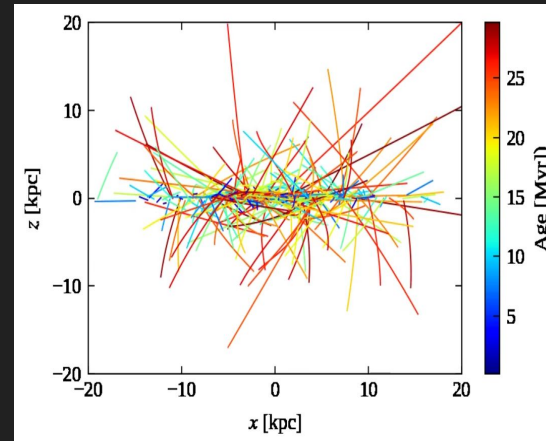
- For our Galactic model  $\Phi_{\text{MW}}$ , we evolve the stars' position & velocity by **solving Newtonian equations of motion** in cylindrical galactocentric coordinates:

$$\ddot{\vec{r}} = -\vec{\nabla}\Phi_{\text{MW}}$$

Top view



Side view

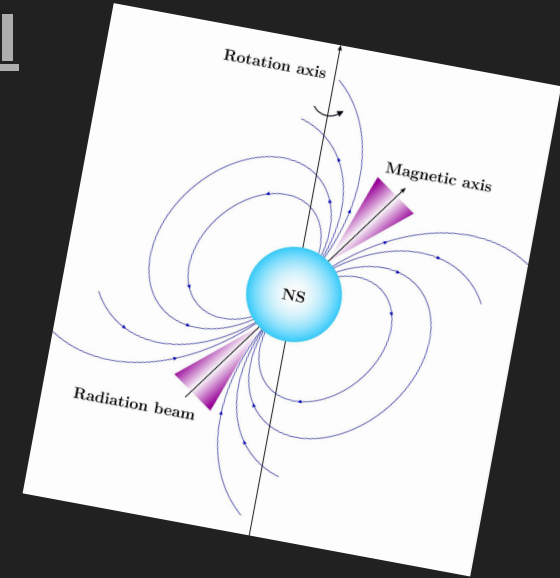


Galactic evolution tracks for  $h_c = 0.18$  kpc,  $\sigma = 265$  km/s.



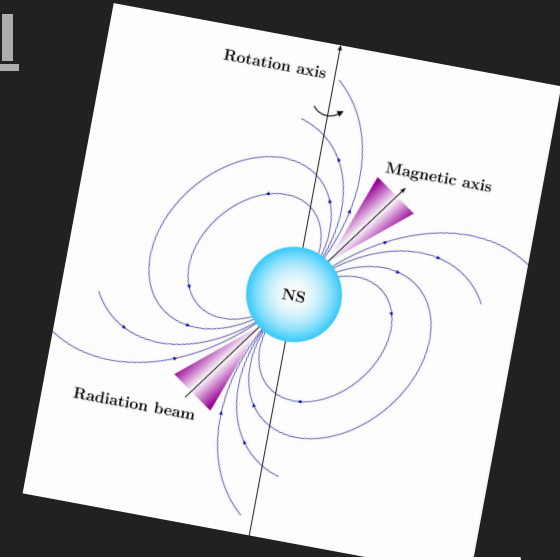
# Magneto-rotational evolution I

- The neutron-star magnetosphere exerts a **torque onto the star**. This causes **spin-down** and the **alignment of the magnetic and rotation axes**.
- Neutron star **magnetic fields decay** due to the Hall effect and Ohmic dissipation in the outer stellar layer (crust) (e.g., Viganó et al., 2013 & 2021).
- We make the following assumptions:
  - **Initial periods** follow a log-normal with  $\mu_{\log P}$  and  $\sigma_{\log P}$  (Igoshev et al., 2022)
  - **Initial fields** follow a log-normal with  $\mu_{\log B}$  and  $\sigma_{\log B}$  (Gullón et al., 2014)
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$$\mathcal{P}(\log P_0) = \frac{1}{\sqrt{2\pi}\sigma_{\log P}} \exp\left(-\frac{\log P_0 - \mu_{\log P}}{2\sigma_{\log P}^2}\right)$$

**We vary the five free parameters  $\mu_{\log P}$ ,  $\mu_{\log B}$ ,  $\sigma_{\log P}$ ,  $\sigma_{\log B}$ , and  $a$ .**

# Magneto-rotational evolution II

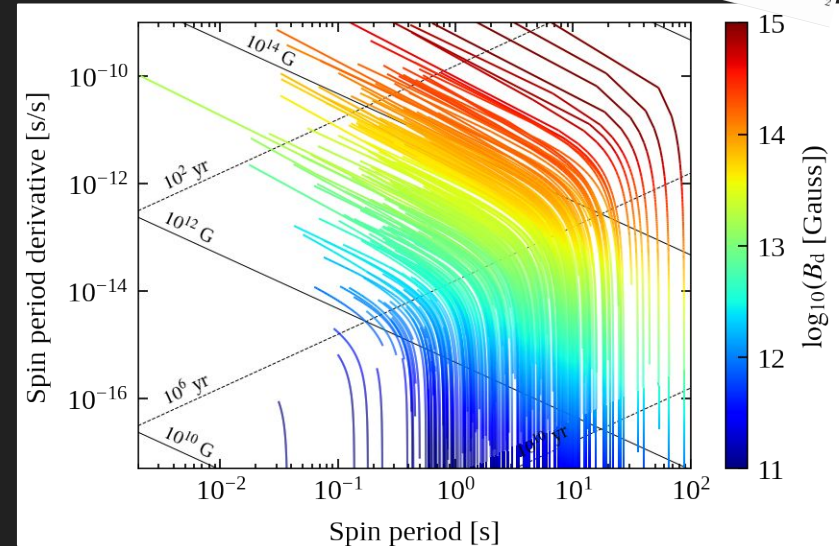
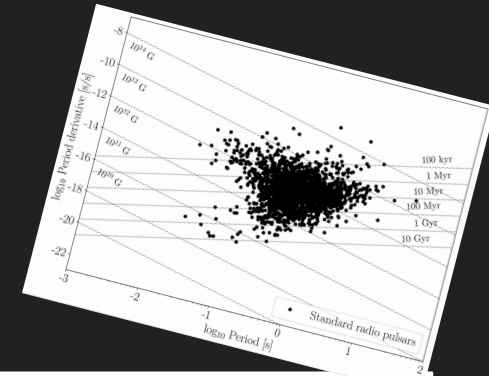
- To model the magneto-rotational evolution, we numerically **solve two coupled ordinary differential equations** for the period and the misalignment angle (Aguilera et al., 2008; Philippov et al. 2014).
- We use results from **2D magneto-thermal simulations** to determine the evolution of the magnetic field.
- This allows us to follow the stars'  $P$  and  $\dot{P}$  evolution in the  $PP$ -plane.

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Period period-derivative evolution tracks for

$$\mu_{\log P} = -0.6, \sigma_{\log P} = 0.3, \\ \mu_{\log B} = 13.25, \sigma_{\log B} = 0.75.$$



# Radio emission and detection

- The stars' **rotational energy**  $E_{\text{rot}}$  is converted into coherent radio emission. We assume that the corresponding **radio luminosity**  $L_{\text{radio}}$  is proportional to the loss of  $E_{\text{rot}}$  (Faucher- Giguère & Kaspi, 2006; Gullón et al., 2014).  $L_0$  is taken from observations.

$$L_{\text{radio}} = L_0 \left( \frac{\dot{P}}{P^3} \right)^{1/2} \propto \dot{E}_{\text{dot}}^{1/2}$$

- As the radio **emission is beamed**, between ~60-90% of pulsars do not point towards us. For those that do intercept our line of sight, we compute the **radio flux**  $S_{\text{radio}}$  and the **pulse width**  $W$ .

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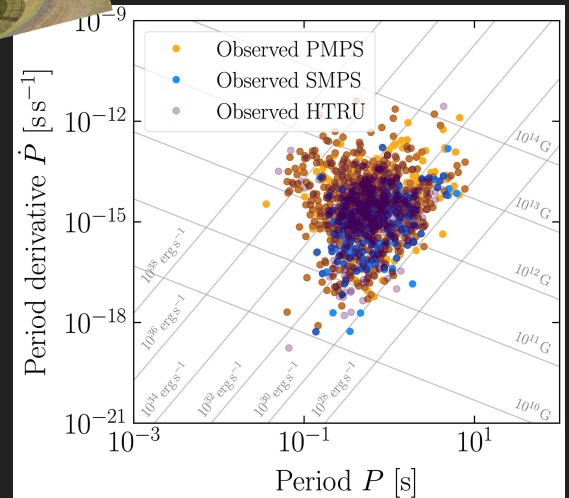
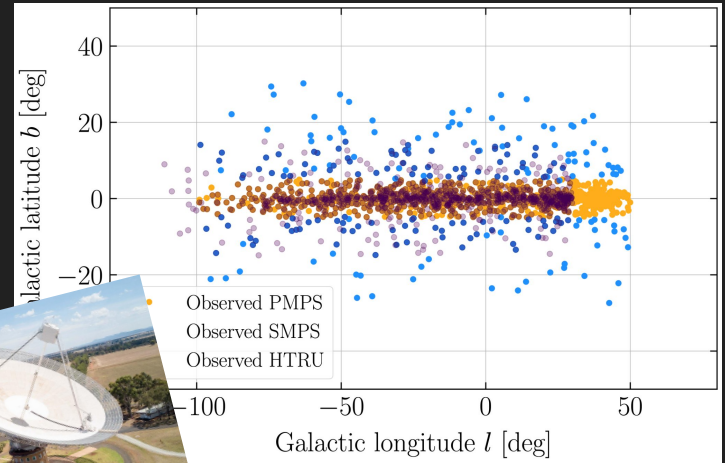
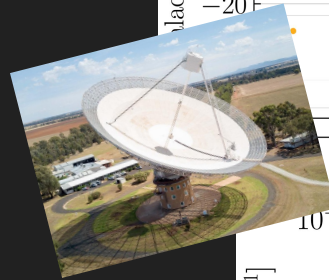
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A pulsar counts as detected, if it **exceeds the sensitivity threshold** for a survey recorded with a specific radio telescope.

# Three pulsar surveys

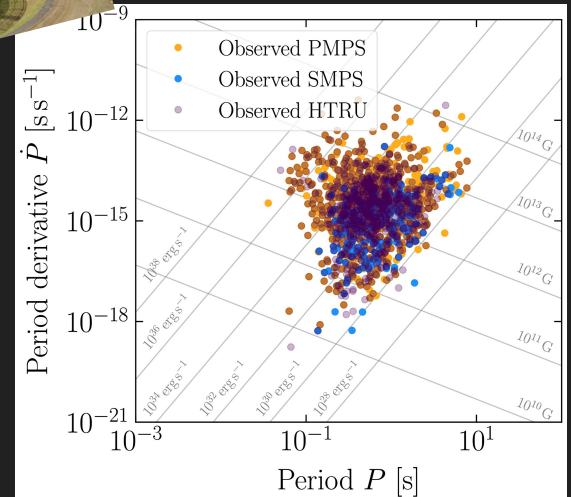
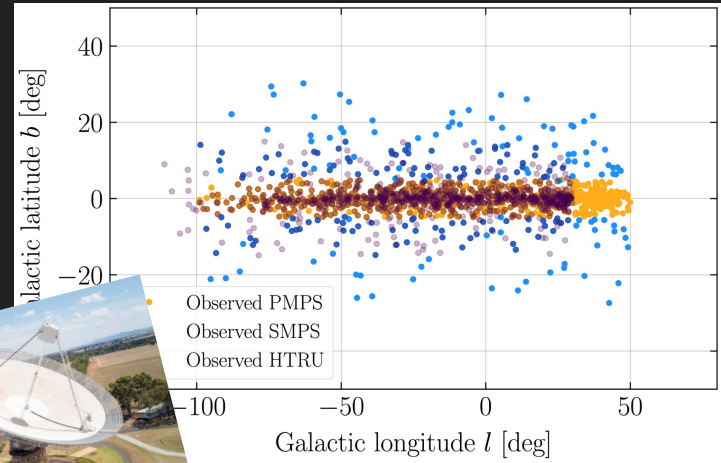
- We compare our simulated populations with three surveys from Murriyang (the Parkes Radio Telescope):
  - **Parkes Multibeam Pulsar Survey (PMPS):** 1,009 isolated pulsars
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Can we constrain birth properties by looking at a current snapshot of the pulsar population?





# Comparing models and data

- Population synthesis models are complex and have **many free parameters**. Traditional methods (e.g., by-eye studies,  $\chi^2$ - and KS-tests, annealing methods, MCMC) for parameter inference scale poorly and require simplified models.
- To constrain our free parameters, we want to perform **Bayesian inference**: based on some prior knowledge  $\pi(\theta)$ , a stochastic model and an observation  $x$ , infer the most likely distribution  $P(\theta|x)$  for our model parameters  $\theta$  given:

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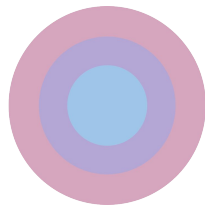
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For complex simulators, the **likelihood is defined implicitly and often intractable**. This is overcome with **simulation-based** (likelihood-free) **inference** (see e.g. Cranmer et al., 2020).

# Simulation-based inference

- To perform **Bayesian inference for any kind of (stochastic) forward model** (e.g. those specified by simulators), we can use the following approach:

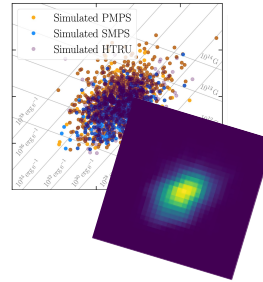
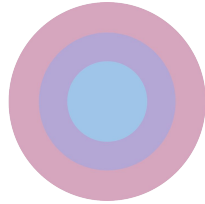
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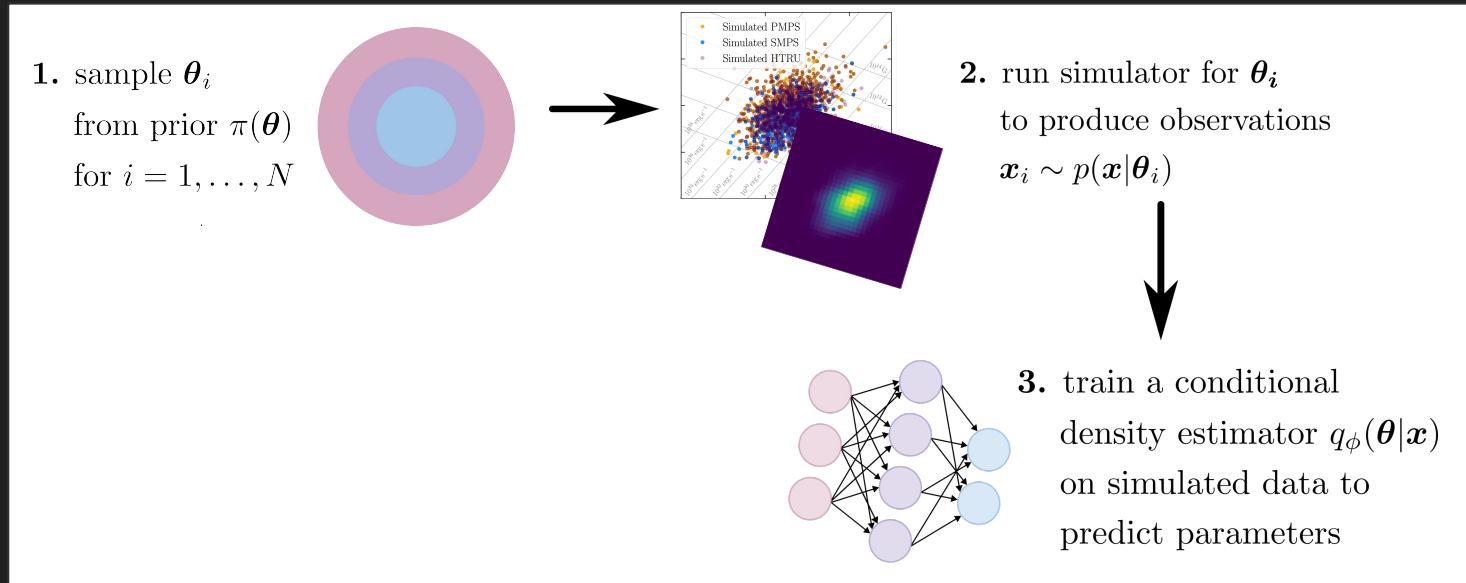
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2. run simulator for  $\theta_i$   
to produce observations  
 $\mathbf{x}_i \sim p(\mathbf{x}|\theta_i)$

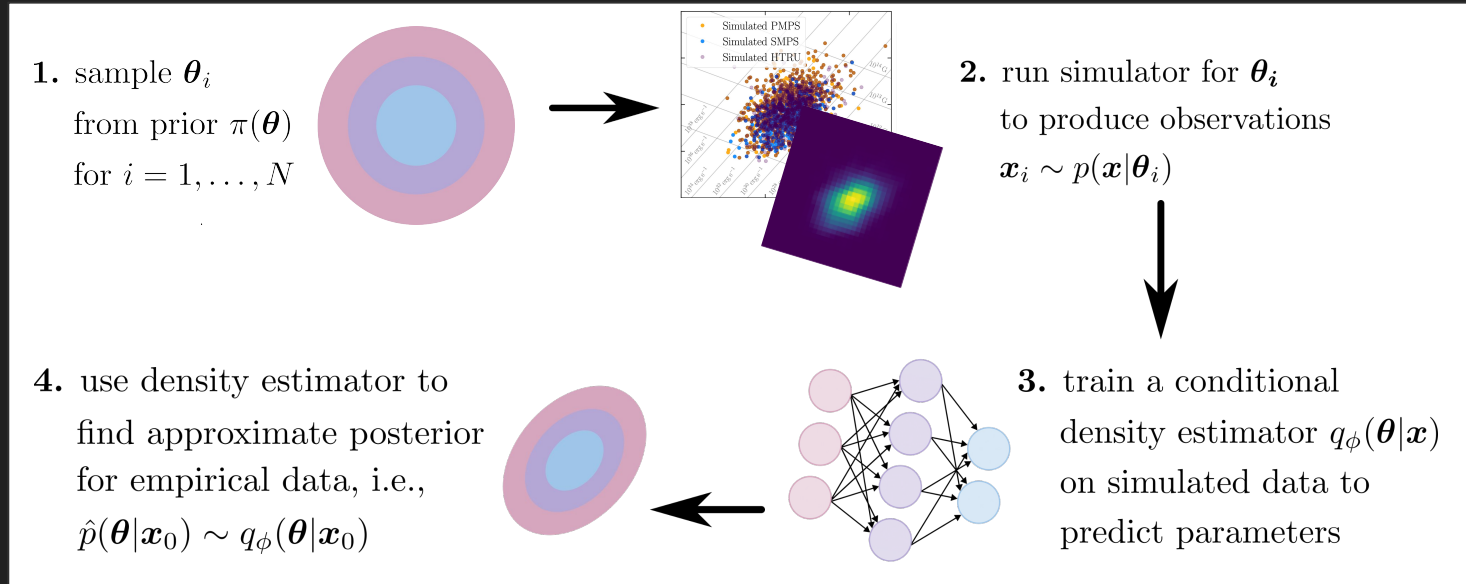
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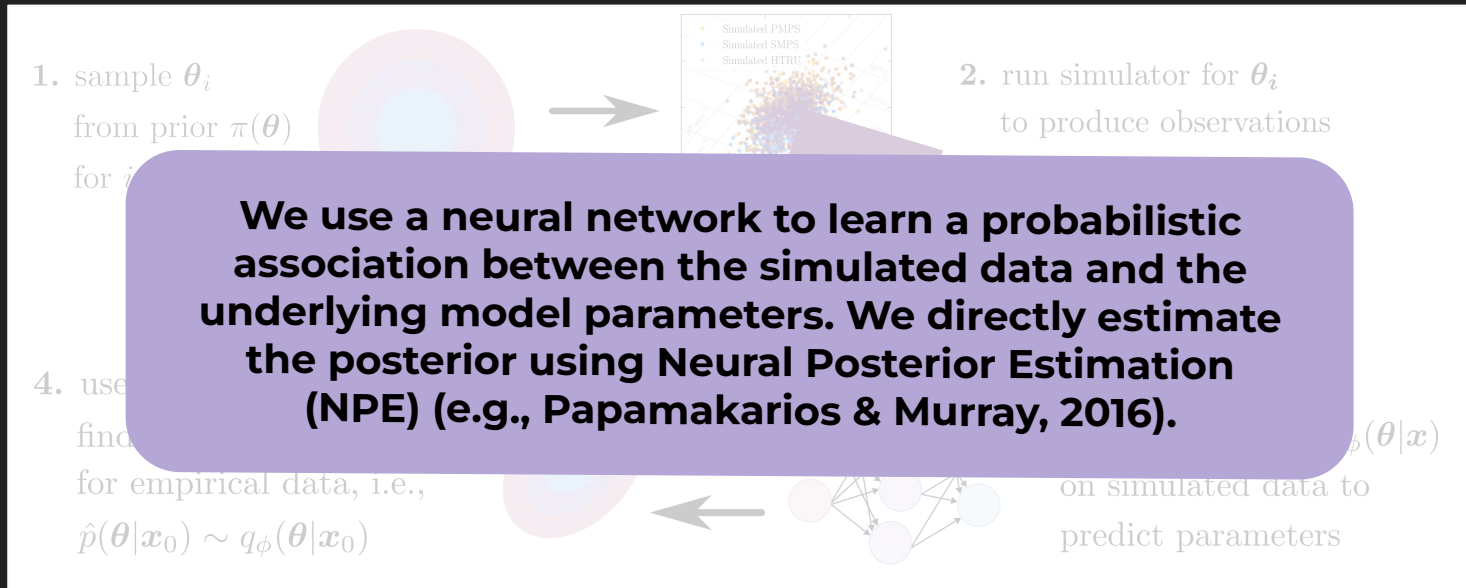
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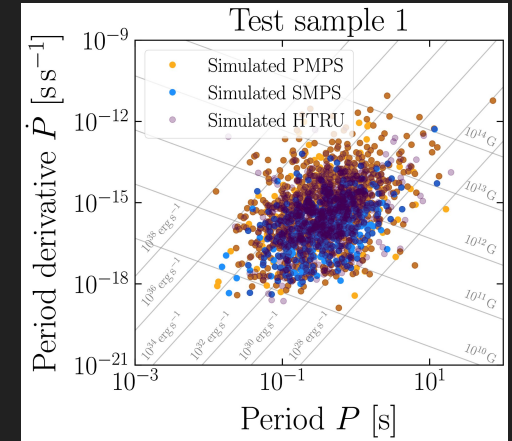


# Simulator output

- We simulate 360,000 populations by varying the 5 magneto-rotational parameters and generate density maps for 3 surveys in  $P\dot{P}$ -plane (Graber et al., 2024).

- $\mu_{\log P} \in \text{Uniform}(-1.5, -0.3)$
- $\sigma_{\log P} \in \text{Uniform}(0.1, 1.0)$
- $\mu_{\log B} \in \text{Uniform}(12.0, 14.0)$
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Test simulation for  
 $\mu_{\log P} = -0.85, \sigma_{\log P} = 0.51,$   
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and  $a_{\text{late}} = -0.86.$



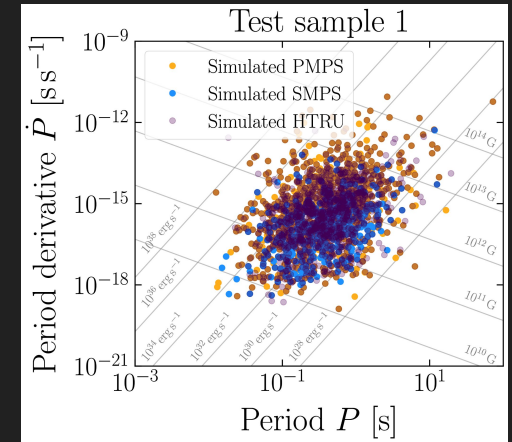
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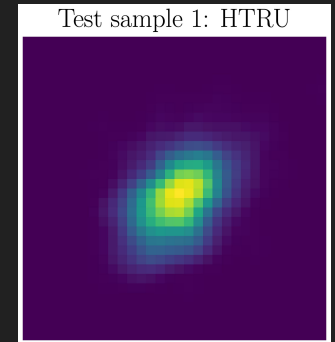
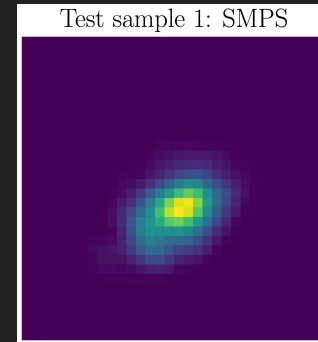
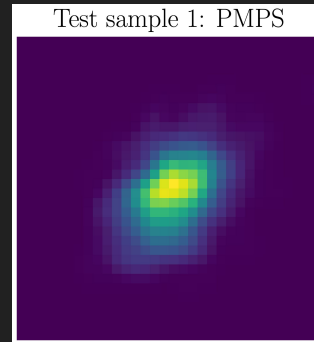
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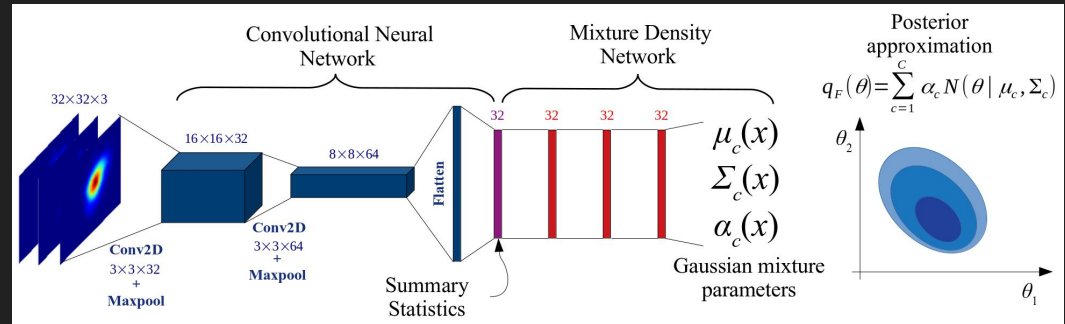


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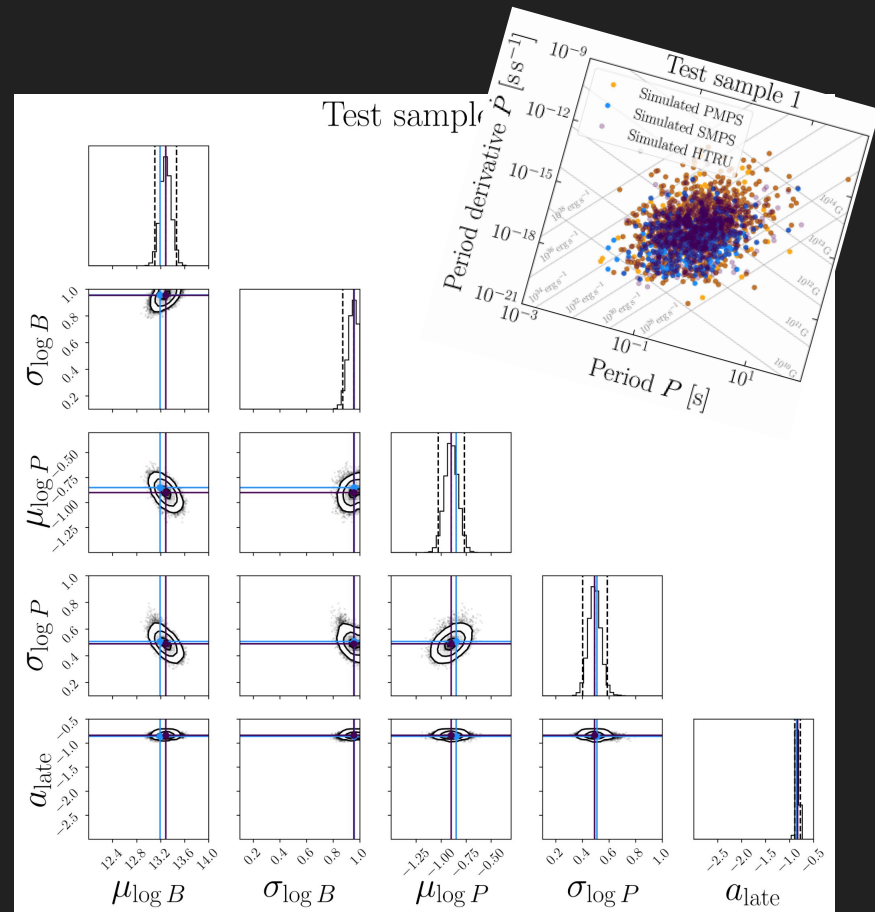
# Neural network setup

- To perform SBI, we use the **PyTorch package sbi** (Tejero-Cantero et al., 2020; <https://www.mackelab.org/sbi/>). Our trainable neural network has two parts:
  - **CNN** (see Ronchi et al., 2021): compresses the maps into a latent vector.
  - **Mixture density network (MDN)**: posterior is approximated by a mixture of 10 Gaussians components; we learn the means, stds and coefficients.
- We use Kaiming initialization for the CNN and Xavier for the MDN, 89% of data for training, 10% for validation and 1% for testing, a batch size of 8 & learning rate of  $5 \times 10^{-4}$ .



# Inference on test sample

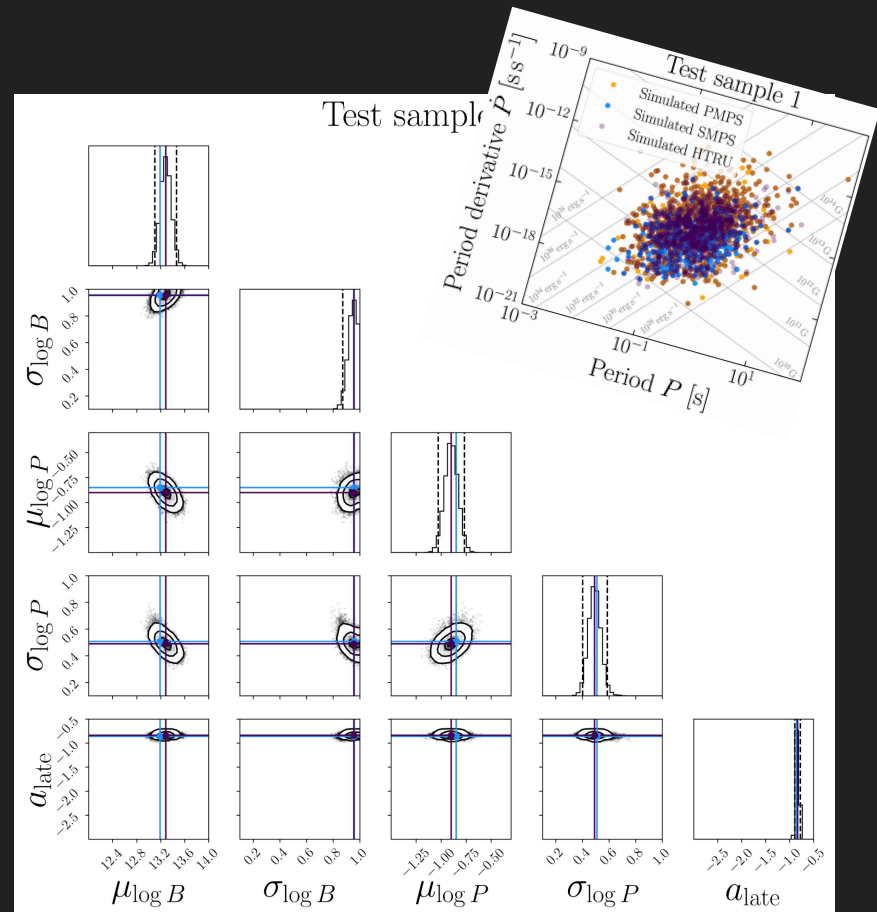
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We recover **narrow and well-defined posteriors** for all five parameters that typically contain the ground truth (parameters used for the forward simulation) at the 95% credibility level.

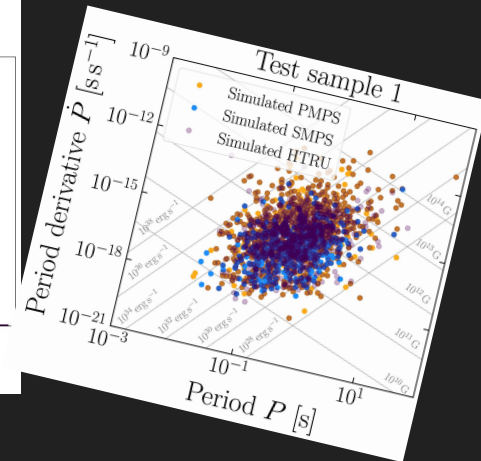
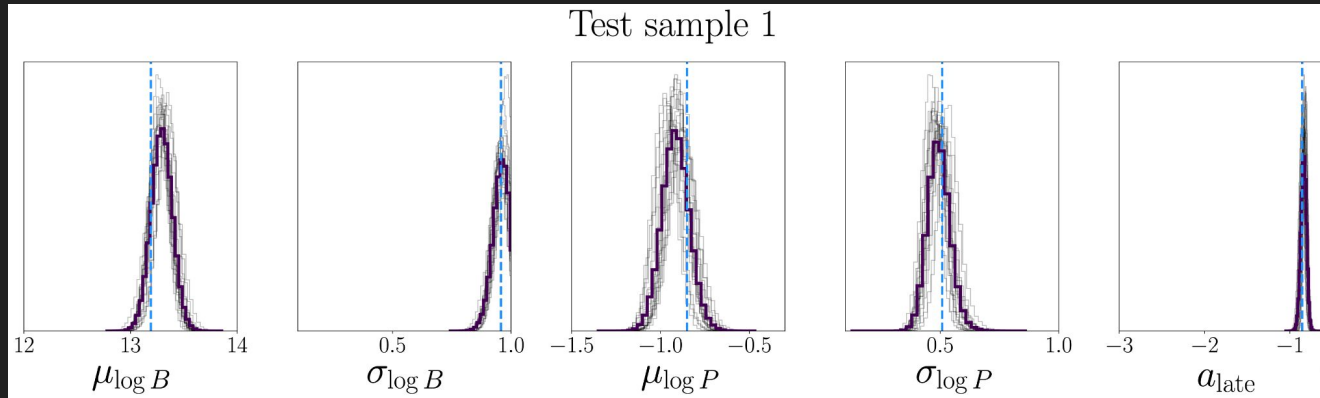


# Ensemble of networks

- Our inferred posteriors are **sensitive to neural network setup**. To analyse the robustness of our results, we perform **22 ML experiments** by varying network hyper parameters, survey input, map resolution, preprocessing.
- We combine posteriors from 19 experiments with well-behaved training behaviour into an **ensemble posterior to obtain conservative predictions**.

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# Observed population

- We apply our ensemble network to infer on the observed population of isolated radio pulsars at 95% CI:

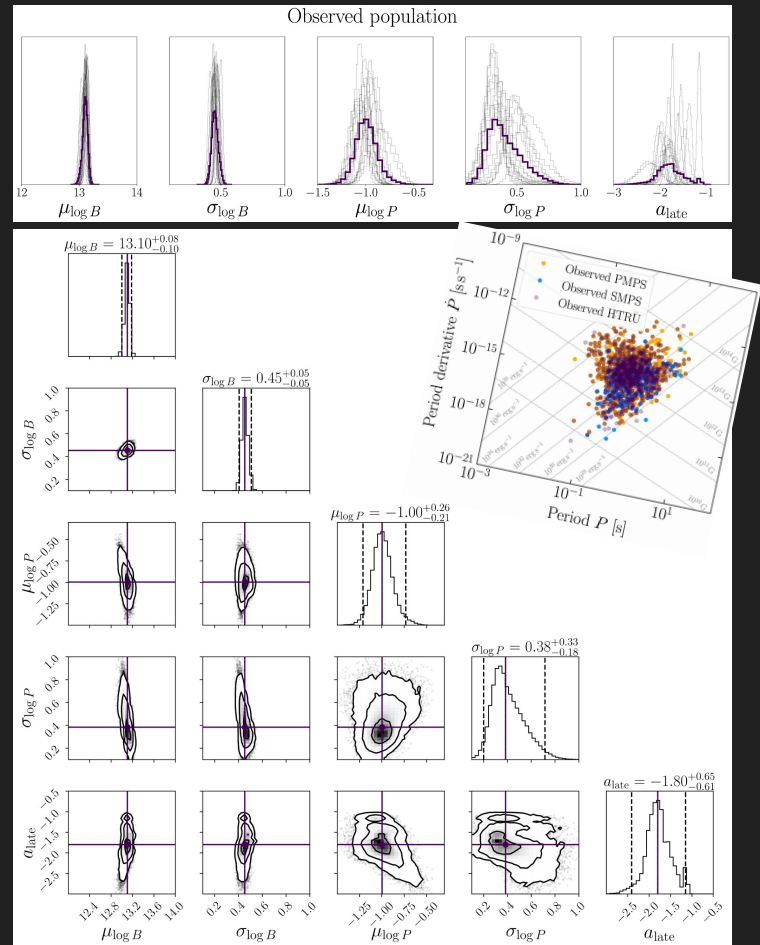
$$\mu_{\log B} = 13.10^{+0.08}_{-0.10},$$

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$$\sigma_{\log P} = 0.38^{+0.33}_{-0.18},$$

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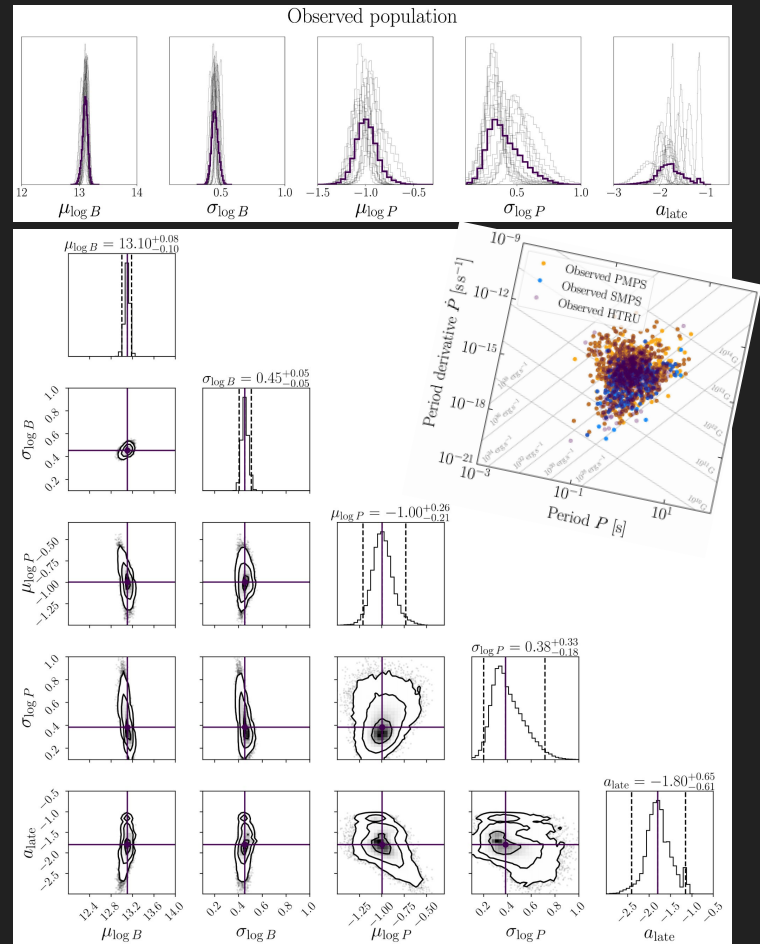
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$$\sigma_{\log P} = 0.38_{-0.18}^{+0.33}$$

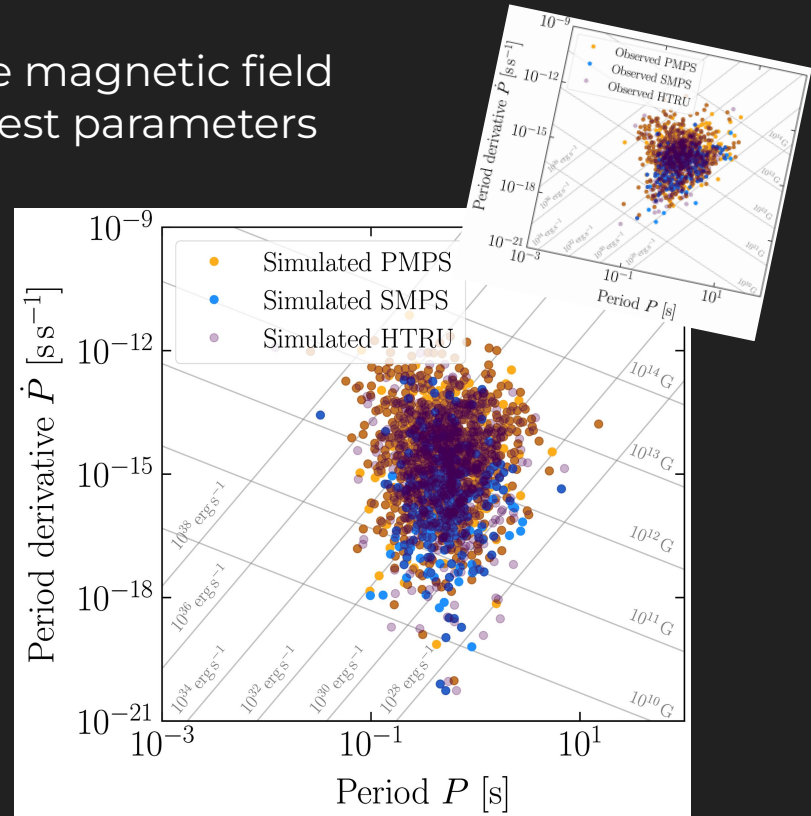
$$a_{\text{late}} = -1.80_{-0.61}^{+0.65}$$

We find **tight constraints for  $B_0$  and  $P_0$**  parameters. Larger uncertainties and bi-modal posterior in  $a_{\text{late}}$  hint at insufficient modelling of late-time B-field decay.



# Two sanity checks

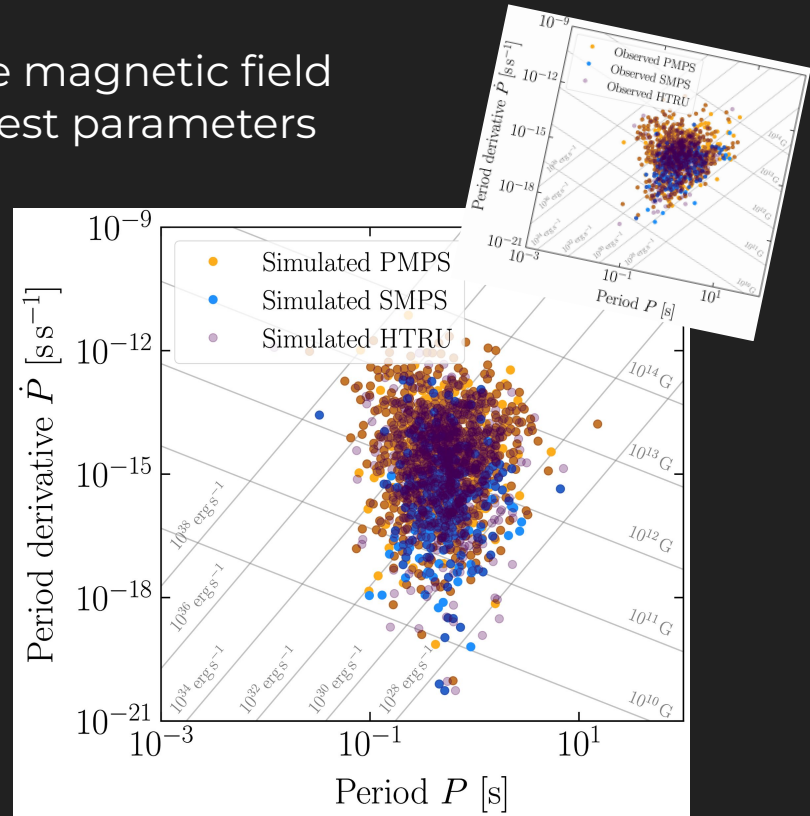
- Despite potential issues with the late-time magnetic field decay, rerunning our simulator with the best parameters produces a **reasonable PP diagram**.
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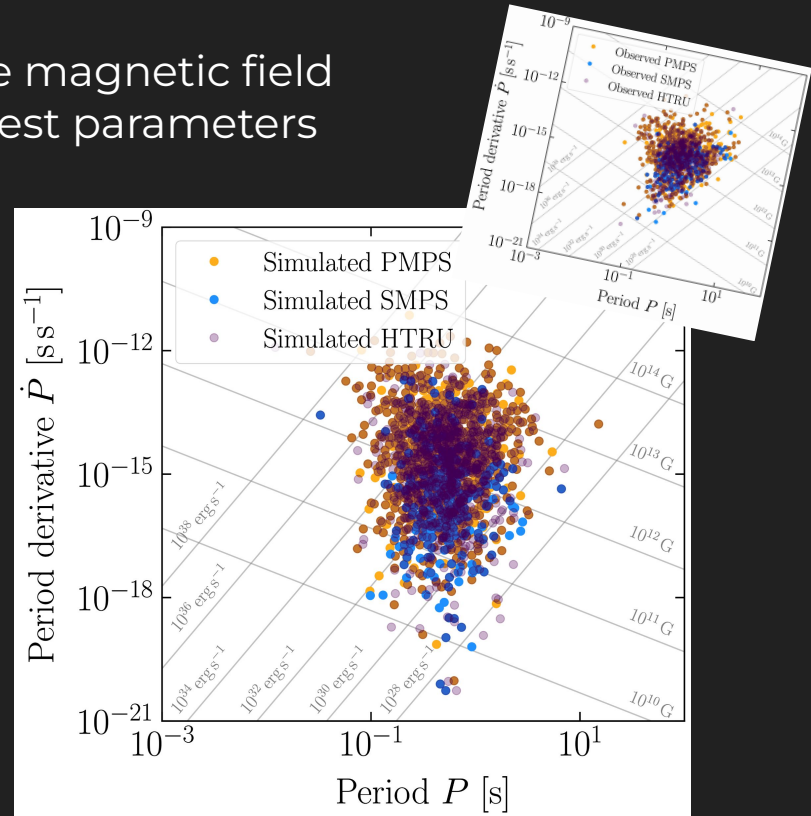


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This is in good agreement with core-collapse supernova rates, i.e.,  $1.63 \pm 0.46$  per century (Rozwadowska et al., 2021).



# Take-home points and outlook

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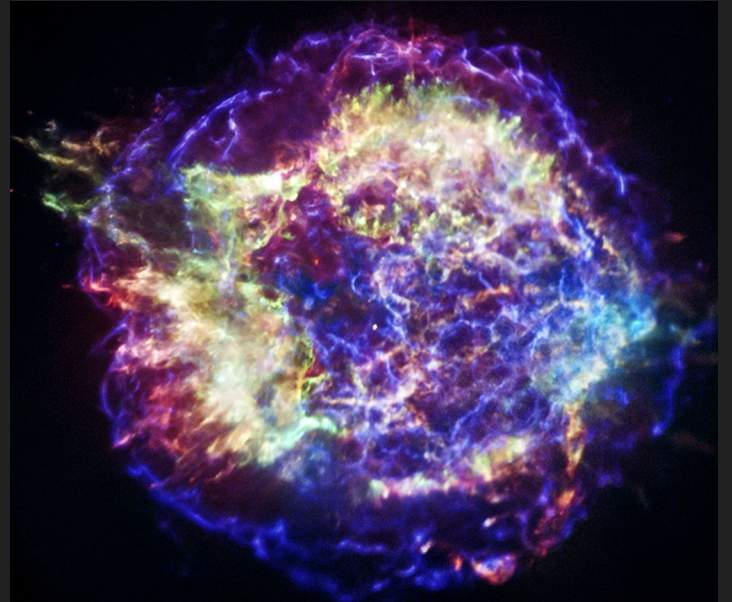
- Population synthesis bridges gap between detected pulsars and the invisible population.
- It allows us to constrain birth rates and NS birth properties.

- SBI has opened up the possibility for robust statistical inference with complex simulators.
- We successfully used SBI to infer magneto-rotational parameters.

- We are currently working to improve our approach:
  - Model multi-wavelength emission and observations.
  - Remove the need for large training datasets by using a sequential SBI approach that adaptively simulates data.



**THANK YOU**



Cassiopeia A supernova remnant  
(credit: NASA/CXC/SAO)



# Magneto-rotational evolution - analytics

- For a given initial rotation period, magnetic field strength and inclination angle (uniformly distributed along the sphere), the evolution of a neutron star is determined by **three ordinary differential equations**. The first two are

$$\dot{P} = \frac{\pi^2 B^2 R^6}{c^3 IP} (\kappa_0 + \kappa_1 \sin^2 \chi)$$

$$\dot{\chi} = -\frac{\pi^2 B^2 R^6}{c^3 IP^2} (\kappa_2 \sin \chi \cos \chi)$$

The  $\kappa$  are determined from simulations. For a **realistic pulsar magnetosphere** filled with plasma we have  $\kappa_0 \approx \kappa_1 \approx \kappa_2 \approx 1$  (Spitkovsky 2006; Philippov et al., 2014).

- From the induction equation, we deduce for the B-field (Aguilera et al., 2008):

$$\dot{B} = -\frac{c^2 B}{4\pi\sigma L^2} - \frac{cB^2}{4\pi en_e L^2} \quad \text{with} \quad \tau_{\text{Ohm}} = \frac{4\pi\sigma L^2}{c^2} \sim 10^6 \text{ yr}, \quad \tau_{\text{Hall}} = \frac{4\pi en_e L^2}{cB} \sim 10^4 B_{12} \text{ yr}$$

# Improved B-field prescription

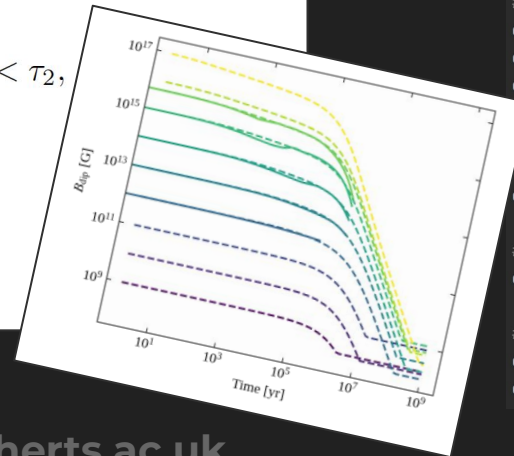
- To provide a more **realistic description of the magnetic field decay**, we **fit an analytical function** to B-field evolution curves obtained from 2D magneto-thermal simulations (Viganó et al., 2021) for various fields. As they are only valid up to  $10^6$  yr, we extend the **late-time behaviour with a power law** as follows:

$$B(t) = B_0 \left(1 + \frac{t}{\tau_1}\right)^{a_1} \left(1 + \frac{t}{\tau_2}\right)^{a_2 - a_1} \left(1 + \frac{t}{\tau_{\text{late}}}\right)^{a_{\text{late}} - a_2} \quad \text{for } \tau_1 < \tau_2 < \tau_{\text{late}},$$

$$B(t) = B_0 \left(1 + \frac{t}{\tau_1}\right)^{a_1} \left(1 + \frac{t}{\tau_{\text{late}}}\right)^{a_{\text{late}} - a_1} \quad \text{for } \tau_1 < \tau_{\text{late}} < \tau_2,$$

$$B(t) = B_0 \left(1 + \frac{t}{\tau_{\text{late}}}\right)^{a_{\text{late}}} \quad \text{for } \tau_{\text{late}} < \tau_1 < \tau_2.$$

with  $\tau_1 \equiv A_1 B_0^{b_1}$   $\tau_2 \equiv A_2 B_0^{b_2}$



```
# Power-law indices.
cfg["a1"]: float = -0.13
cfg["a2"]: float = -3.0

# Timescale parameters, normalization
cfg["A1"]: float = 1.0e14
cfg["b1"]: float = -0.8
cfg["A2"]: float = 6.0e8
cfg["b2"]: float = -0.2

# Timescale in [yr] when transitioning
cfg["tau_late"]: float = 2.0e6

# Late time power-law index.
cfg["a_late"]: float = -2.0

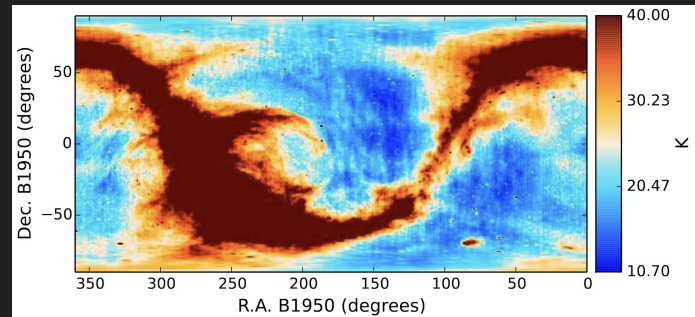
# Parameters for a log-normal distribution
cfg["B_millisec_mean"] = 8.5
cfg["B_millisec_sigma"] = 0.5
```

# Radiometer equation

$$S_{\text{radio}} = \frac{L_{\text{radio}}}{\Omega_{\text{beam}} d^2}$$

- Following Lorimer & Kramer (2005), we convert the intrinsic (bolometric) radio flux  $S_{\text{radio}}$  into a flux density measured at a given frequency. We then account for the fact that the intrinsic pulse width is broadened to  $w_{\text{eff}}$  (due to interstellar scattering, dispersion and finite sampling resolution of a detector) and compute the period-averaged flux  $S_{\text{mean}}$  observed by a radio telescope.
- With these estimates and information specific to the PMPS, SMPS and HTRU surveys, we **determine the signal-to-noise-ratio** for a given simulated pulsar using the **radiometer equation**:

$$S/N = \frac{S_{\text{mean}} G \sqrt{N_{\text{pol}} \Delta\nu \Delta t_{\text{obs}}}}{\beta (T_{\text{rec}} + T_{\text{sky}}(l, b))} \sqrt{\frac{P - w_{\text{eff}}}{w_{\text{eff}}}}$$



Haslam sky temperature map taken from Remazeilles et al., (2015).

# SBI approaches

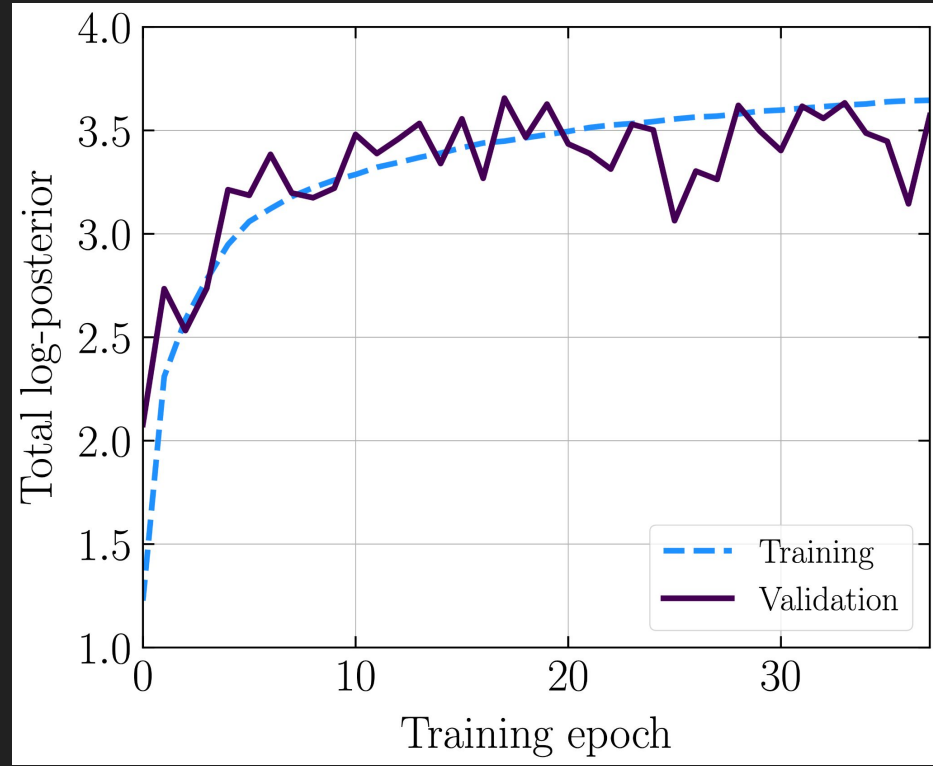
- Different approaches (all relying on deep learning) exist to **learn a probabilistic association** between the simulated data and the underlying parameters. These algorithms essentially focus on different pieces of Bayes' theorem:
  - Neural Posterior Estimation (NPE) (e.g., Papamakarios & Murray, 2016)
  - Neural Likelihood Estimation (NLE) (e.g., Papamakarios et al., 2019)
  - Neural Ratio Estimation (NRE) (e.g., Hermans et al., 2020; Delaunoy et al., 2022)

**We focus on NPE.** This allows us to **directly learn the posterior distribution**. In contrast, NLE and NRE need an extra (potentially time consuming) MCMC sampling step to construct a posterior.

- All methods exist in **sequential form** (SNPE, SNLE, SNRE), **which adds a fifth step to workflow**. Instead of sampling from the prior, we adaptively generate simulations from the posterior. This **typically requires fewer simulations**.

# Neural network training behaviour

- **Typical learning behaviour** for our simulation-based inference experiments.
- The **log-probability increases** as a function of training epochs for the training and validation loss.
- We see some variation in the validation loss but **little overfitting**.

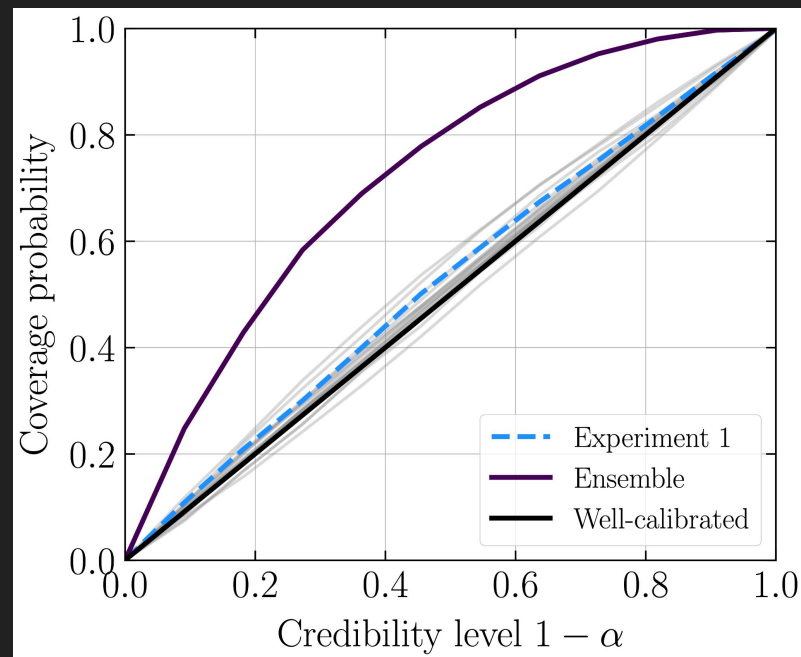


# 22 different ML experiments

No.	Res	Surveys	Frac (%)	Input	Comp	BS	LR	CNN	VM	Epochs	Time (s)	TM
1	32	PMPS, SMPS, HTRU	100	std	10	8	0.0005	baseline	3.65	38	9,373	3.64
2	32	PMPS, SMPS, HTRU	100	std	10	8	0.0005	deep	3.71	49	14,292	3.71
3	64	PMPS, SMPS, HTRU	100	std	10	8	0.0005	baseline	3.55	55	78,837	3.54
4	64	PMPS, SMPS, HTRU	100	std	10	8	0.0005	deep	3.64	89	128,119	3.64
5	32	PMPS, SMPS, HTRU	75	std	10	8	0.0005	baseline	3.74	71	13,232	3.78
6	32	PMPS, SMPS, HTRU	50	std	10	8	0.0005	baseline	3.56	58	7,000	3.55
7*	32	PMPS, SMPS, HTRU	100	norm	10	8	0.01	baseline	3.47	30	7,445	3.73
8	32	PMPS, SMPS, HTRU	100	norm	10	8	0.001	baseline	9.66	54	13,015	9.60
9	32	PMPS, SMPS, HTRU	100	std	8	8	0.0005	baseline	3.74	52	12,389	3.73
10	32	PMPS, SMPS, HTRU	100	std	5	8	0.0005	baseline	3.83	118	27,973	3.86
11	32	PMPS, SMPS, HTRU	100	std	10	16	0.0005	baseline	3.99	85	10,476	3.97
12	32	PMPS, SMPS, HTRU	100	std	10	32	0.0005	baseline	4.11	79	5,346	4.06
13	32	PMPS, SMPS, HTRU	100	std	10	8	0.001	baseline	3.36	61	14,785	3.33
14	32	PMPS, SMPS, HTRU	100	std	10	8	0.0001	baseline	4.22	75	18,369	4.22
15	32	HTRU	100	std	10	8	0.0005	baseline	3.43	63	15,568	3.42
16	32	SMPS, HTRU	100	std	10	8	0.0005	baseline	3.58	40	9,979	3.59
17	32	PMPS, SMPS	100	std	10	8	0.0005	baseline	3.41	69	16,937	3.41
18*	64	PMPS, SMPS, HTRU	50	std	10	8	0.0005	baseline	3.45	47	5,766	3.44
19	32	PMPS, SMPS, HTRU	100	norm	10	32	0.001	baseline	10.05	44	2,864	10.20
20	32	PMPS, SMPS, HTRU	100	norm	10	32	0.0001	baseline	10.31	90	5,815	10.49
21	32	PMPS, SMPS, HTRU	100	norm	10	16	0.001	baseline	9.82	77	9,901	9.98
22*	32	PMPS, SMPS, HTRU	100	norm	10	16	0.0001	baseline	10.45	124	15,603	10.55

# Quality check: coverage calculation

- A measure that provides information on the quality of our posterior is the **coverage**. I.e., we take a frequentist approach and determine how often our true parameters fall into the  $(1-\alpha)\%$  credible interval under the model.
- By definition, a  **$(1-\alpha)\%$  credibility level covers the true parameter in  $(1-\alpha)\%$  of our analyses**, resulting in a diagonal line. In this case, the **posterior is neither conservative (too wide; above diagonal) nor overconfident (too narrow; below diagonal)**.



# Comparison to other studies

- A comparison with earlier studies is difficult due to different simulation models and inference methods.

References	$\mu_{\log B}$	$\sigma_{\log B}$	$\mu_{\log P}$	$\sigma_{\log P}$
Faucher-Giguère & Kaspi (2006)	12.65	0.55	—	—
Gullón et al. (2015)	12.99	0.56	—	—
Ciešlar et al. (2020)	$12.67^{+0.01}_{-0.02}$	$0.34^{+0.02}_{-0.01}$	—	—
Igoshev et al. (2022)	12.44	0.44	$-1.04^{+0.15}_{-0.20}$	$0.53^{+0.12}_{-0.08}$
This work	$13.10^{+0.04}_{-0.05}$	$0.45^{+0.03}_{-0.02}$	$-1.00^{+0.11}_{-0.10}$	$0.38^{+0.16}_{-0.10}$

